

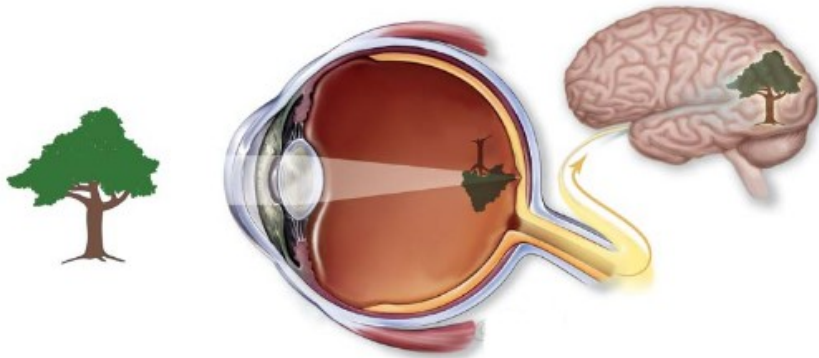
A universal framework for quantum theory from qubits to quantum gravity

Robert Oeckl

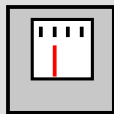
Centro de Ciencias Matemáticas
Universidad Nacional Autónoma de México
Morelia, Mexico

Gravity in Qubits
Bratislava, Slovakia
23 November 2018

Vision and reconstruction



<http://tpe.vision.aveugles.free.fr/vision.php>



Operational approach:
Fundamental notions

- experiment
- measurement
- observation
- preparation
- intervention

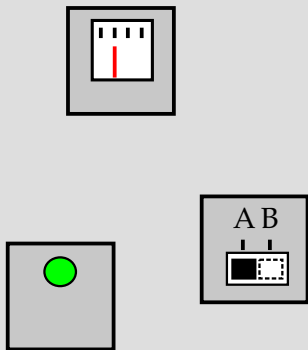
Subsume instance as:

- **process**

Processes have
outcomes.

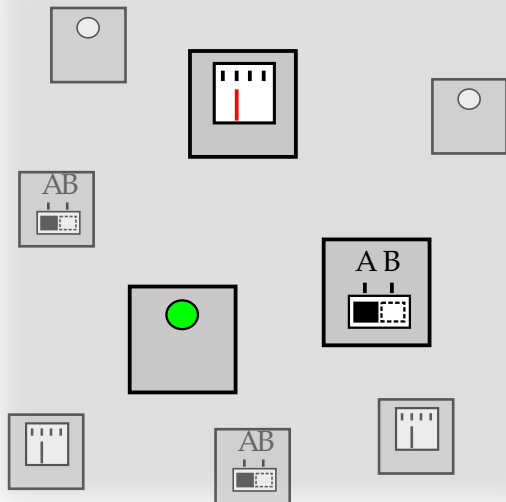
Represent processes as
boxes.

Processes and interfaces



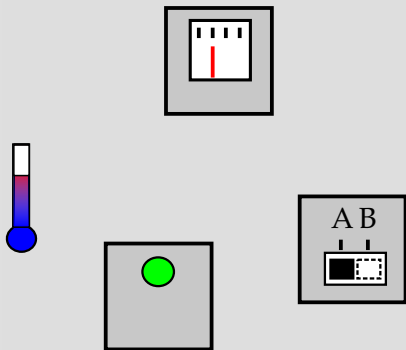
Processes are not isolated. Outcomes depend on other processes. We want to predict **correlations**.

Processes and interfaces



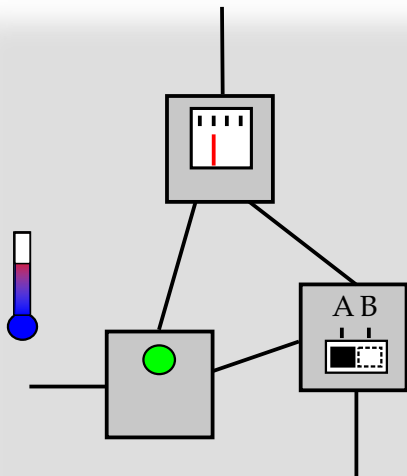
The outcome of a given set of processes depends generally on a large number of other processes.

Processes and interfaces



We treat these external processes collectively. We call this a **boundary condition**.

Processes and interfaces

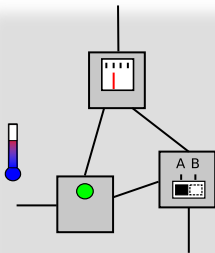


We introduce the notion of **interface** to model **interaction** between processes. An interface encodes **communication** or **information exchange** between processes. we depict this as a **link**.

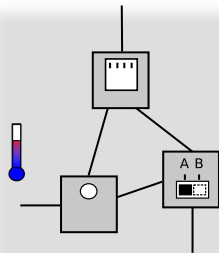
Open links carry boundary conditions.

Processes and interfaces form **networks**, represented as **labeled graphs**.

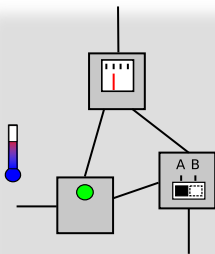
Probabilities



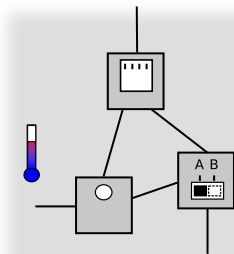
To obtain a
probability we
need to ask for
← this
conditional on
this →



Probabilities



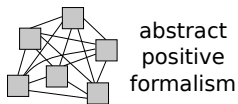
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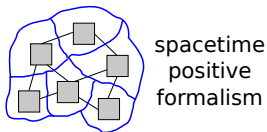
Diagrams translate to **positive real numbers**. Probability is a **quotient**.

$$P = \frac{\begin{array}{|c|} \hline \text{Green dot} \\ \hline \end{array} \begin{array}{|c|} \hline \text{A B} \\ \hline \end{array} \begin{array}{|c|} \hline \text{Scale} \\ \hline \end{array} \begin{array}{|c|} \hline \text{Thermometer} \\ \hline \end{array}}{\begin{array}{|c|} \hline \text{White dot} \\ \hline \end{array} \begin{array}{|c|} \hline \text{A B} \\ \hline \end{array} \begin{array}{|c|} \hline \text{Scale} \\ \hline \end{array} \begin{array}{|c|} \hline \text{Thermometer} \\ \hline \end{array}}$$

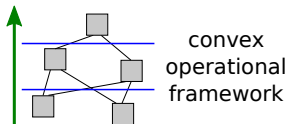
$$P = \frac{P(g) \diamond Q(A) \diamond R([x_1, x_2]) \diamond b}{P(*) \diamond Q(A) \diamond R(*) \diamond b}$$



+ spacetime + locality



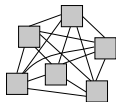
+ time + causality



classical
(lattices)

quantum
(anti-lattices)

abstract
classical
statistical
theory



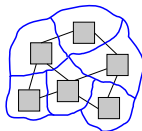
abstract
positive
formalism

abstract
quantum
theory



+ spacetime + locality

spacetime
statistical
field theory



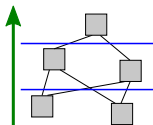
spacetime
positive
formalism

general
boundary
formulation
/
axiomatic
QFT



+ time + causality

statistical
mechanics



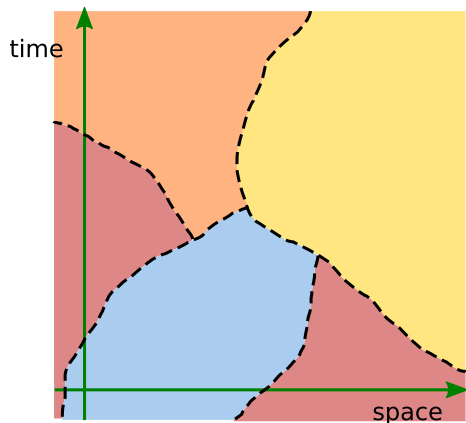
convex
operational
framework

standard
formulation
of quantum
theory

The local positive formalism

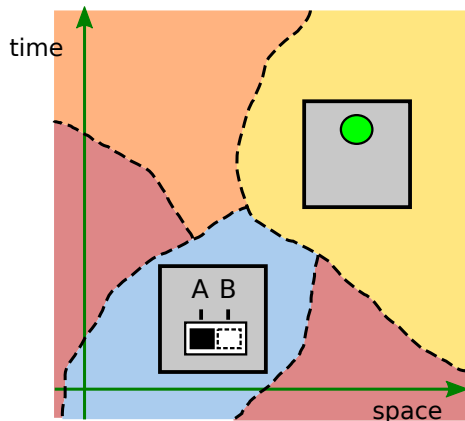
Spacetime locality provides a powerful organizing principle. Processes only interface with **adjacent** processes. This decreases considerably the inter-connectivity of the graph.

The local positive formalism



For a **local** description cut up spacetime into pieces, called **regions**. These are in contact with each other through **hypersurfaces**.

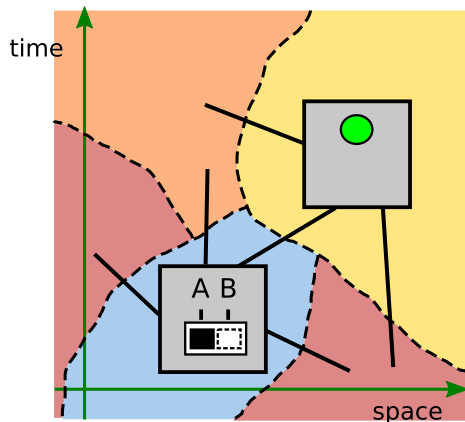
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Associate **processes** to regions to encode experiments, observations, interventions etc.

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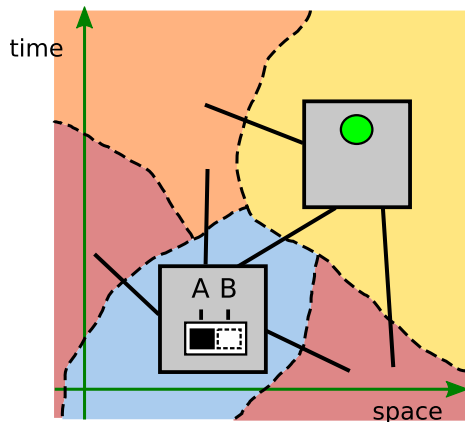


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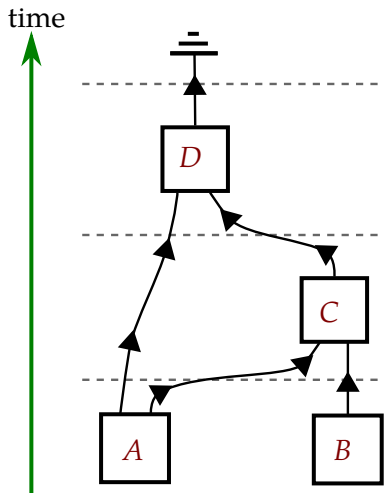
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Associate **processes** to regions to encode experiments, observations, interventions etc.

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Glue regions along hypersurfaces \longleftrightarrow Compose processes along links

Time-evolution positive formalism



Suppose that we have a **fixed notion of time**. The arrangement of processes in time then yields a **partial order**.

We allow links only between processes that are comparable. Links carry **arrows** directed to the future. We obtain a **directed acyclic graph**.

Ingredients

\mathcal{B} real **partially ordered vector space** with inner product (\cdot, \cdot) and **order unit** $\mathbf{e} \in \mathcal{B}$.

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In classical theory: **statistical distributions**

- \mathcal{B} space of square-integrable functions on probability space (L, μ)
- $\langle \sigma, \tau \rangle = \int \sigma(x)\tau(x)d\mu(x)$
- $\mathbf{e} = 1$ the unit constant function on L

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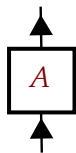
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In quantum theory: **density matrices**

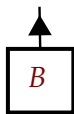
- \mathcal{B} space of self-adjoint operators on the Hilbert space \mathcal{H}
- $\langle \sigma, \tau \rangle = \text{tr}(\sigma\tau)$
- $\mathbf{e} = \text{id}$ the identity operator on \mathcal{H}

Elementary diagrams

Diagrams represent linear maps.



$A : \mathcal{B} \rightarrow \mathcal{B}$
operation



$B : \mathbb{R} \rightarrow \mathcal{B}$
preparation

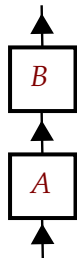


$C : \mathcal{B} \rightarrow \mathbb{R}$
effect

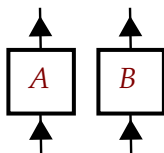


$\sigma \mapsto (\mathbf{e}, \sigma)$
discard

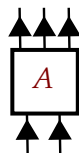
Composition



$$B \circ A : \mathcal{B} \rightarrow \mathcal{B}$$



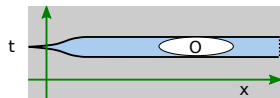
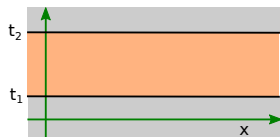
$$A \otimes B : \\ \mathcal{B}_1 \otimes \mathcal{B}_2 \rightarrow \mathcal{B}_1 \otimes \mathcal{B}_2$$



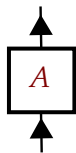
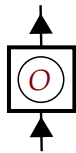
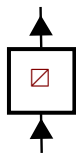
$$A : \mathcal{B}_1 \otimes \mathcal{B}_2 \\ \rightarrow \mathcal{B}_3 \otimes \mathcal{B}_4 \otimes \mathcal{B}_5$$

Classical operations

spacetime



diag.



phase space

$U : L \rightarrow L$
preserve
measure

$O : L \rightarrow \mathbb{R}$

(stat.) state space

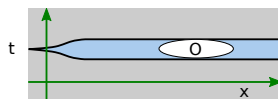
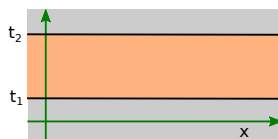
$\hat{U} : \mathcal{B} \rightarrow \mathcal{B}$
 $(\hat{U}(\sigma))(x) = \sigma(U^{-1}(x))$

$\hat{O} : \mathcal{B} \rightarrow \mathcal{B}$
 $(\hat{O}(\sigma))(x) = O(x)\sigma(x)$

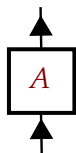
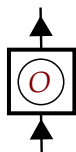
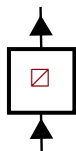
$A : \mathcal{B} \rightarrow \mathcal{B}$
positive

Quantum operations

spacetime



diag.



Hilbert space

$U : \mathcal{H} \rightarrow \mathcal{H}$
unitary

$O : \mathcal{H} \rightarrow \mathcal{H}$
hermitian
 $O = \sum_j o_j P_j$

(stat.) state space

$\hat{U} : \mathcal{B} \rightarrow \mathcal{B}$
 $\hat{U}(\sigma) = U\sigma U^\dagger$

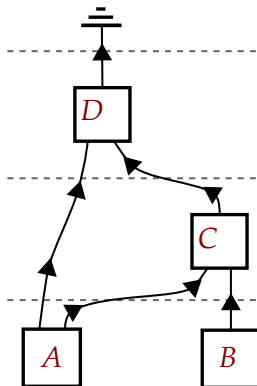
$\hat{O} : \mathcal{B} \rightarrow \mathcal{B}$
 $\hat{O}[i](\sigma) = P_i \sigma P_i^\dagger$
 $\hat{O}[*](\sigma) = \sum_j \hat{O}[j]$
 $\hat{O}(\sigma) = \sum_j o_j \hat{O}[j]$

$A : \mathcal{B} \rightarrow \mathcal{B}$
completely positive
 $A(\sigma) = \sum_j K_j \sigma K_j^\dagger$

Evaluation and probability

Example:

- preparations A, B
- measurements C, D
- a discard effect e

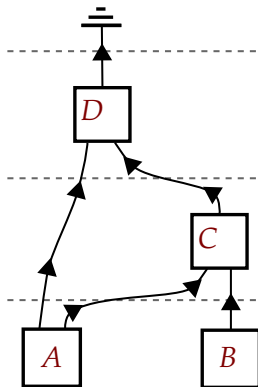


Evaluation and probability

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What is the probability for outcome (i, j) in measurements (C, D) ?

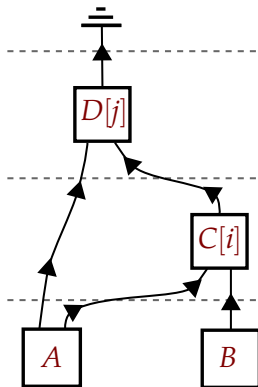


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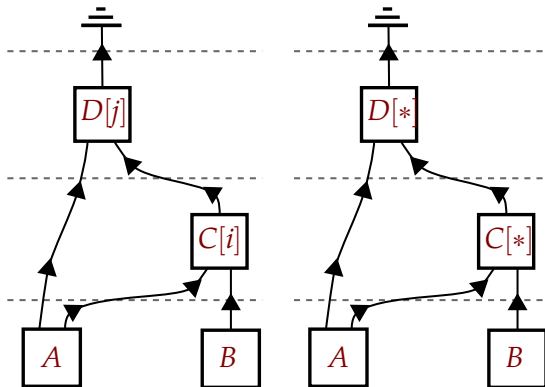


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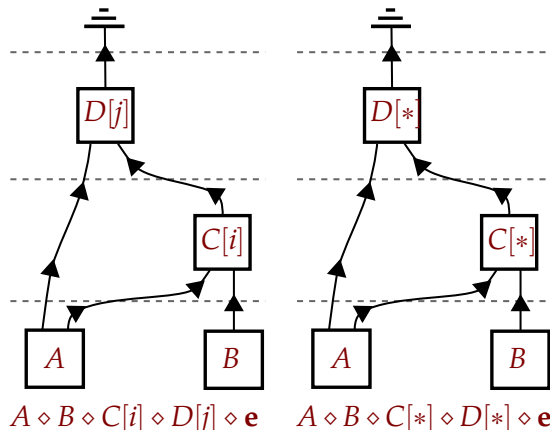


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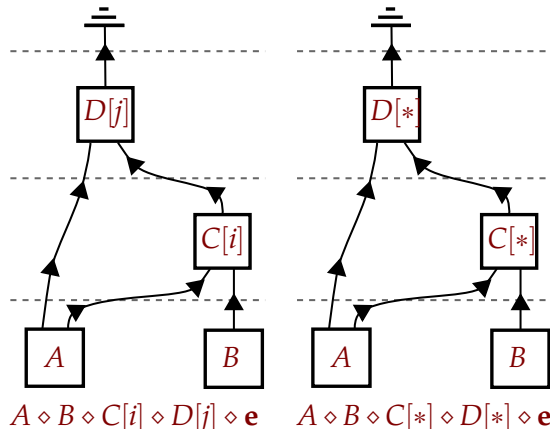


Evaluation and probability

Example:

- preparations A, B
- measurements C, D
- a discard effect e

What is the probability for outcome (i, j) in measurements (C, D) ?



Probability:
$$P(i, j) = \frac{A \diamond B \diamond C[i] \diamond D[j] \diamond e}{A \diamond B \diamond C[*] \diamond D[*] \diamond e}$$

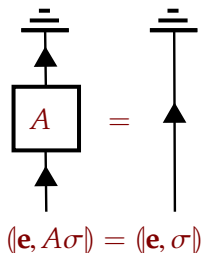
Causality

Causality is implemented through a normalization condition on operations. This is **time-asymmetric**, **requires directionality**.

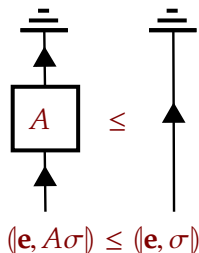
Causality

Causality is implemented through a normalization condition on operations. This is **time-asymmetric**, **requires directionality**.

$A : \mathcal{B} \rightarrow \mathcal{B}$ non-selective



$A : \mathcal{B} \rightarrow \mathcal{B}$ selective

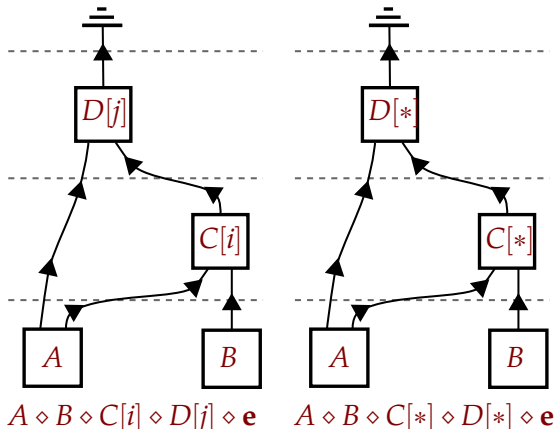


Evaluation and probability (II)

Example:

- preparations A, B
- measurements C, D
- a discard effect \mathbf{e}

What is the probability for outcome (i, j) in measurements (C, D) ?



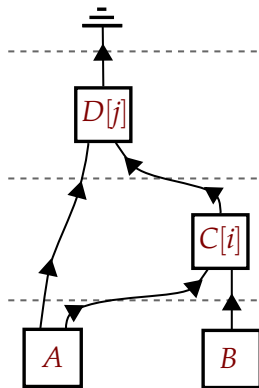
Probability:
$$P(i, j) = \frac{A \diamond B \diamond C[i] \diamond D[j] \diamond \mathbf{e}}{A \diamond B \diamond C[*] \diamond D[*] \diamond \mathbf{e}}$$

Evaluation and probability (II)

Example:

- preparations A, B
- measurements C, D
- a discard effect \mathbf{e}

What is the probability for outcome (i, j) in measurements (C, D) ?



Causality implies:

Composites of non-selective operations (including preparations and discard) evaluate to 1.

$$A \diamond B \diamond C[i] \diamond D[j] \diamond \mathbf{e}$$

$$A \diamond B \diamond C[*] \diamond D[*] \diamond \mathbf{e} = 1$$

Probability: $P(i, j) = A \diamond B \diamond C[i] \diamond D[j] \diamond \mathbf{e}$

A lesson for QG:

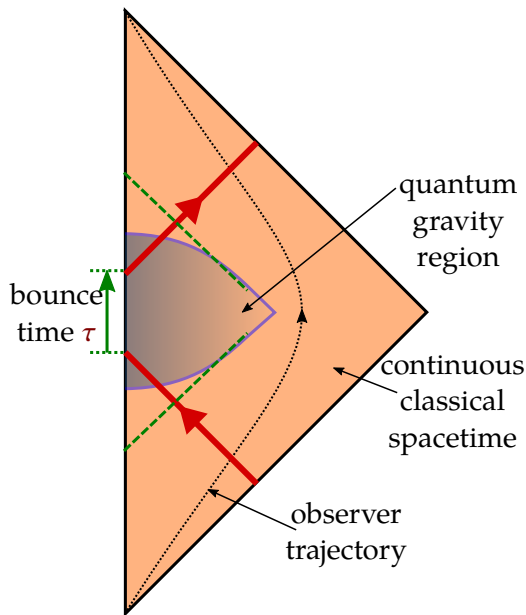
Absent a fixed notion of time, there is no way to consistently implement the standard causality condition. Hence, diagrams **do not** in general evaluate to probabilities, only **quotients** do.

A lesson for QG:

Absent a fixed notion of time, there is no way to consistently implement the standard causality condition. Hence, diagrams **do not** in general evaluate to probabilities, only **quotients** do.

This also applies to the modulus squares of amplitudes!

Example: Black hole bounce

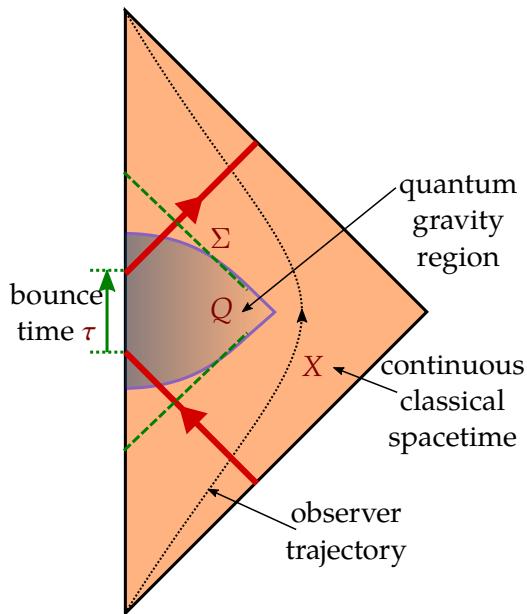


Black hole to white hole bounce model.

Two parameters:

- shell mass m
- bounce time τ

Example: Black hole bounce



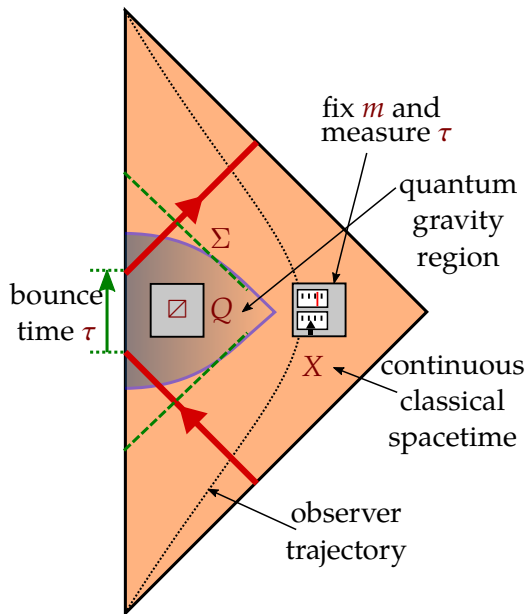
Black hole to white hole bounce model.

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Distinguish spacetime regions Q and X , with interfacing hypersurface Σ .

Example: Black hole bounce



Black hole to white hole bounce model.

Two parameters:

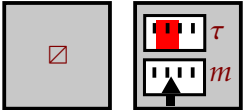
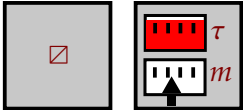
- shell mass m
- bounce time τ

Distinguish spacetime regions Q and X , with interfacing hypersurface Σ .

Introduce two processes with one link.

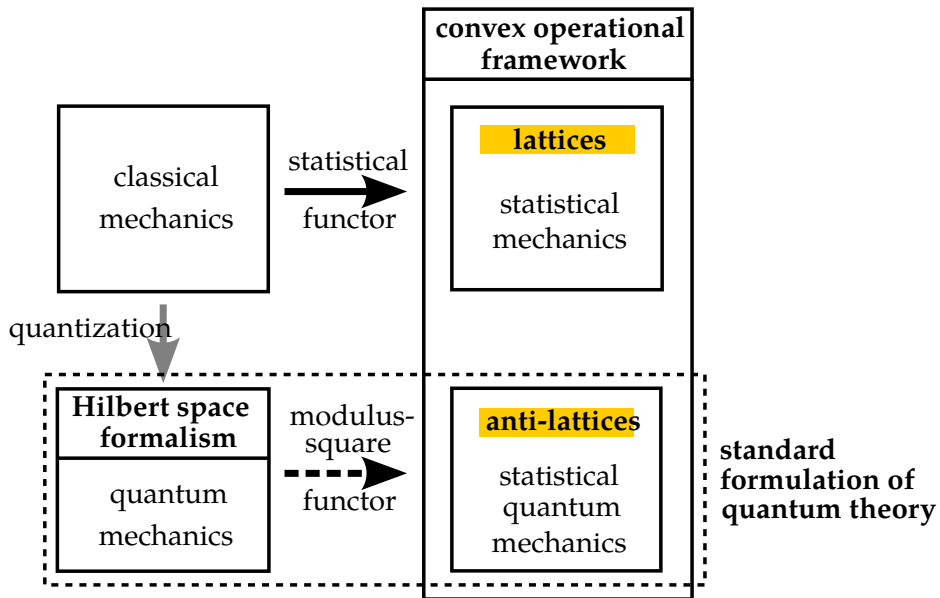
Example: Black hole bounce

Given a fixed shell mass m , we wish to determine the **probability** for the bounce time τ to lie in a given interval $[\tau_1, \tau_2]$. This is,

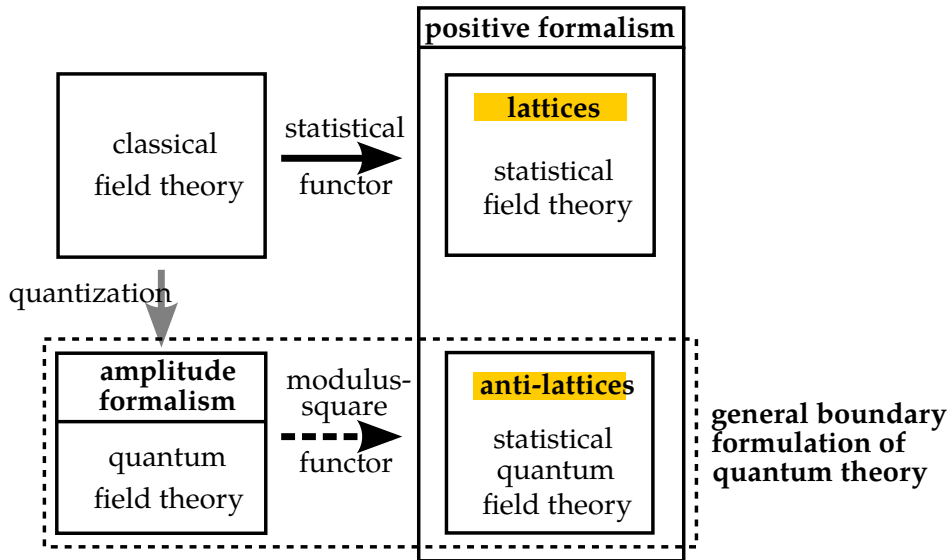
$$P(m, [\tau_1, \tau_2]) = \frac{\text{Diagram 1}}{\text{Diagram 2}}$$

$$P(m, [\tau_1, \tau_2]) = \frac{A_Q \diamond A_X[m, [\tau_1, \tau_2]]}{A_Q \diamond A_X[m, [0, \infty])}$$


R. O., *A predictive framework for quantum gravity and black hole to white hole transition*, Phys. Lett. A **82** (2018) 2622–2625, arXiv:1804.02428.

Fundamental physics: Time-evolution frameworks



Fundamental physics: Spacetime frameworks



Some applications in QFT

- Non-linear models:
 - ▶ **Three dimensional quantum gravity** is a TQFT and fits “automatically”. [Witten 1988;...]
 - ▶ **Quantum Yang-Mills theory** in 2 dimensions for arbitrary regions and hypersurfaces with corners. [RO 2006]
 - ▶ **abelian Yang-Mills theory** in higher dimensions. [Díaz 2014–; Díaz, RO 2017]
- New **S-matrix** type asymptotic amplitudes [Colosi, RO 2008; Colosi 2009; Dohse 2011; 2012]
- QFT in **curved spacetime**: dS, AdS and more [Colosi, Dohse 2009–]
- **Rigorous and functorial quantization** of linear and affine field theories without metric background. [RO 2010; 2011; 2012]
- **Unruh effect**. [Colosi, Rätzel 2012; Bianchi, Haggard, Rovelli 2013]
- Striking results for **fermions**: Hilbert spaces become **Krein spaces** and an **emergent notion of time**. [RO 2012]
- Partial solution of **state-locality** in fermionic QFT [RO 2013]

Conclusions

- The positive formalism **extends the scope of quantum theory** to situations where no metric background is present. It is thus a suitable **basis for formulating quantum theories of gravity**.
- **Loop quantum gravity** already heavily uses the (pure state version) of the positive formalism.
- The positive formalism **unifies** the essentials of **various formalisms** that have been developed for quantum foundations in recent years (causaloid formalism, quantum combs, operator-tensor formalism, process matrix, categorical approaches, ...)
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→ The **positive formalism / GBF** appears ideally suited as a basic ingredient for diverse contributions to the QSS project.

Main reference

R. O., *A local and operational framework for the foundations of physics*, to appear in *Adv. Theor. Math. Phys.*, arXiv:1610.09052.