

Quantization of string-like sources  
coupled to BF theory :  
transition amplitudes and topological invariance

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## Outline

- Introduction
- Geometrical interpretation
- Quantization : physical inner product
- Topological invariance

## Introduction

- 2 + 1 gravity : Particle worldline  $\gamma \equiv$  local conical singularity in spacetime curvature

$$F[A] = p \delta_\gamma \quad (1)$$

- How to extend the same idea to higher dimensions ?

$\Rightarrow$  Replace particles by  $(d - 3)$ -branes

(Baez, Wise, Crans - 06; Baez, Perez - 06)

- Framework :  $d$ -dimensional BF theory - compact, semi-simple structure group  $G$

-  $M$  :  $d$ -dimensional spacetime manifold

-  $(A, B)$  :  $\mathfrak{g}$ -valued fields on  $M$  (connection one-form and  $(d - 2)$ -form respectively)

$$S_{BF} = \int_M \text{tr} (B \wedge F[A]) \quad (2)$$

## Introduction

- Idea :

- Fix  $W \subset M \equiv (d - 2)$ -dimensional worldsheet embedded into  $M$
- Put a  $\mathfrak{g}$ -valued  $(d - 3)$ -form  $q$  and a  $G$ -valued function  $\lambda$  on  $W$ , fix a constant unit element  $v \in \mathfrak{g}$  and define the ‘momentum density’  $p = \tau Ad_\lambda(v) \in C^\infty(W, \mathfrak{g})$
- Add the following term to the free action

(de Sousa Gerbert - 90; Baez, Perez - 06)

$$S = \int_M \text{tr} (B \wedge F[A]) - \int_W \text{tr} ((B + d_A q) p), \quad (3)$$

→ Equations of motion

$$\begin{aligned} F[A] &= p \delta_W & , & & d_A B &= [p, q] \delta_W \\ d_A p &= 0 & , & & d_A q &= -B \end{aligned}$$

## Introduction

- Symmetries :  
-standard YM gauge symmetries

$$\begin{aligned}\forall g \in C^\infty(M, G), \quad B &\mapsto B = gBg^{-1} & (4) \\ A &\mapsto A = gAg^{-1} + gdg^{-1} \\ \lambda &\mapsto g\lambda \\ q &\mapsto gqg^{-1}\end{aligned}$$

- ‘topological’ transformations

$$\begin{aligned}\forall \eta \in \Omega^{d-3}(M, \mathfrak{g}), \quad B &\mapsto B + d_A \eta & (5) \\ A &\mapsto A \\ \lambda &\mapsto \lambda \\ q &\mapsto q - \eta\end{aligned}$$

### Geometrical interpretation

In this section,  $d = 4$  and  $G = \text{SO}(\eta)$  ( $\eta = (\pm, +, +, +)$ ) with  $V_\eta = \mathbb{R}\{e_I\}_I$  the **vector representation** space of  $G$

- Two questions :

- What is the **physical meaning** of the algebraic variables  $\lambda$  and  $q$  ?

- Does the theory relate to **solutions to GR** ?

→ Physical interpretation : **Matter on flat backgrounds**

(cf Alejandro's talk)

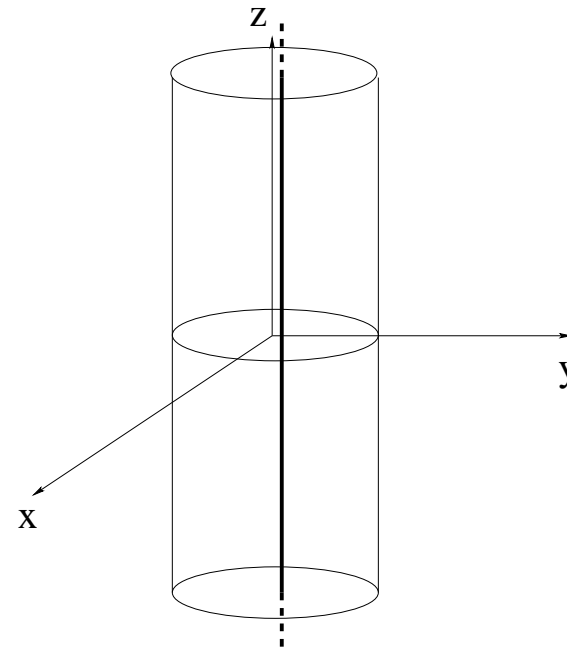
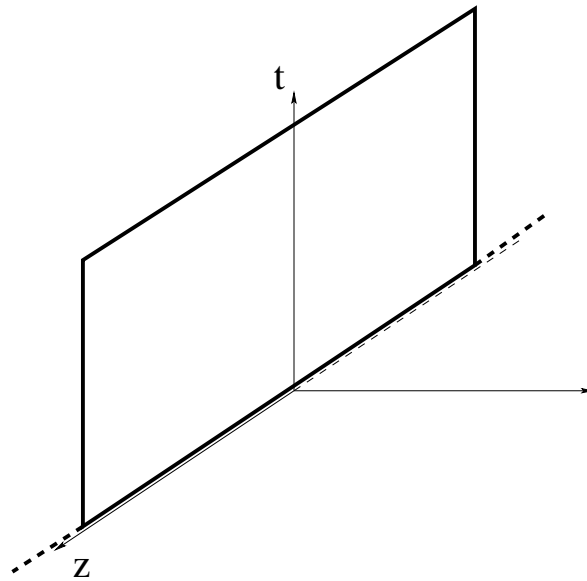
Geometrical interpretation

→ Geometrical interpretation : Cosmic string solutions

- Cosmic string  $\equiv$  infinitely long and thin, straight string

- Let  $\mathcal{S}$  be a spinless cosmic string with mass per unit length  $\tau$

- In local cylindrical coordinates centered on  $\mathcal{S}$  in which the string lies along the  $z$  axis,



Geometrical interpretation

spacetime is described by the dual co-frame  $e^I = e^I_\mu dx^\mu$

$$\begin{aligned}e^0 &= dt \\e^1 &= \cos \varphi dr - 4\tau r \sin \varphi d\varphi \\e^2 &= \sin \varphi dr + 4\tau r \cos \varphi d\varphi \\e^3 &= dz,\end{aligned}\tag{6}$$

and the connection

$$A = A^I_{\mu} \sigma_{IJ} dx^\mu = 4\tau \sigma_{12} d\varphi\tag{7}$$

-The associated spacetime curvature  $F = F^{12} \sigma_{12}$  is singular at the location of the string worldsheet :

$$F^{12} = 8\pi\tau \delta^2(r) dx \wedge dy\tag{8}$$



Geometrical interpretation

-This cosmic string solution can be generated by the Palatini action (+ interaction)

$$S = \int_M \text{tr} ((*e \wedge e) \wedge F[A]) - \tau \int_W \text{tr} ((*e \wedge e) \sigma_{12}) \quad (9)$$

or equivalently by the Plebanski-like action

$$S = \int_M \text{tr} (B \wedge F[A]) - \tau \int_W \text{tr} (Bv) + \frac{1}{2} \int_M \text{tr} (B \wedge \Phi(B)) \quad (10)$$

where we have set  $v = \sigma_{12}$

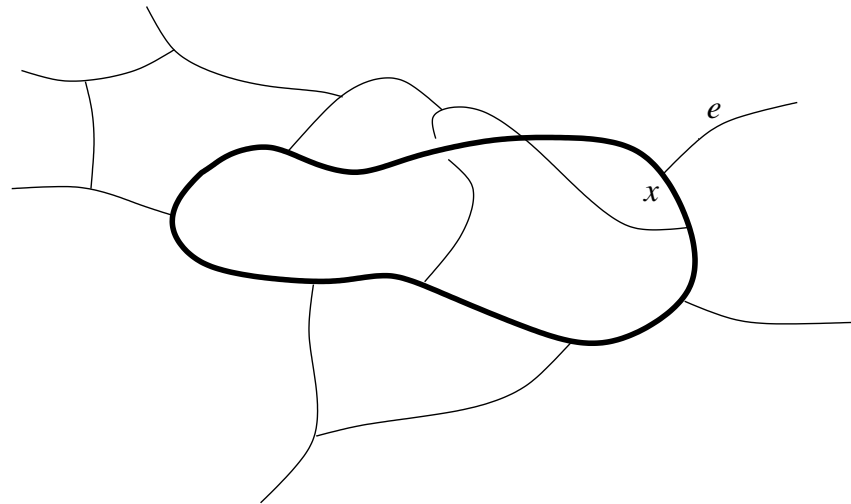
$\Rightarrow$  This is the action of string-like sources coupled to BF theory (in a particular gauge) augmented by the Plebanski term

Quantization : physical inner product

- Canonical analysis :  $M = \Sigma \times \mathbb{R}$  -  $\mathcal{S} = \Sigma \cap W$ 
  - Gauss law :  $G_i = \epsilon^{abc} D_a B_{bci} + \int_{\mathcal{S}} \dot{x}^a [q_a, p]_i \delta_{\mathcal{S}} = 0$
  - Curvature constraint :  $C_i^a = \epsilon^{abc} F_{bci} - \int_{\mathcal{S}} \dot{x}^a p_i \delta_{\mathcal{S}} = 0$
- Canonical quantization of the kinematics:
  - $\mathcal{H}_{kin} = \mathbb{C}\{\Psi_{\alpha}\}_{\alpha}$
  - $\Psi_{\alpha} \equiv$  string spin networks (SSN) (Thiemann - 97; Baez, Perez - 06)
- SSN : Open graph  $\Gamma$  - finite set of points  $X$  on  $\mathcal{S}$ 
  - Edges and endpoints → unitary, irreducible representations  $\rho$  of  $G$
  - Vertices (including endpoints) → intertwining operators  $\iota$

Quantization : physical inner product

$$\Psi_{\Gamma, X}[A, \lambda] = \left[ \bigotimes_{e \in \Gamma} \rho_e[g_e] \bigotimes_{x \in X} \rho_x[\lambda_x] \right] \cdot \bigotimes_{v \in \Gamma} \iota_v, \quad (11)$$



- Kinematical inner product  $\langle \Psi_{\Gamma, X}, \Psi'_{\Gamma, X} \rangle$ : AL-measure  
 $\rightarrow$  Haar integrals assigned to the edges  $(\prod_e \int_G dg_e)$  and endpoints  $(\prod_x \int_G d\lambda_x)$  of  $(\Gamma, X)$

- **Physical inner product** : Formal definition

→ Introduce the **rigging map** (Rovelli, Reisenberger - 97)

$$\eta_{\text{phys}} : \text{Cyl} \rightarrow \text{Cyl}^*; \Psi \mapsto \delta(\hat{C})\Psi$$

$$\Rightarrow \eta_{\text{phys}}(\text{Cyl}) = \text{Cyl}_{\text{phys}}^* \subset \text{Cyl}^*$$

with  $\text{Cyl}_{\text{phys}}^* \equiv$  vector space of **solutions to the curvature constraint**

→ **Physical inner product** :

$$\langle \eta_{\text{phys}}(\Psi_1), \eta_{\text{phys}}(\Psi_2) \rangle_{\text{phys}} = [\eta_{\text{phys}}(\Psi_2)](\Psi_1) \quad (12)$$

$$= \langle \Psi_1, \delta(\hat{C})\Psi_2 \rangle \quad (13)$$

- **Regularization** of the physical inner product

→ make sense of  $\delta(\hat{C})$

$$\delta(\hat{C}) = \prod_{x \in \Sigma} \delta(\hat{C}(x)) = \int_{\mathcal{N}} \mathcal{D}N \exp \left( i \int_{\Sigma} \text{tr}(N \wedge \hat{C}) \right) \quad (14)$$

Quantization : physical inner product

→ Interesting duality :

$$\int_{\Sigma} \text{tr}(N \wedge C) = \int_{\Sigma} \text{tr}(N \wedge F) + \int_{\mathcal{L}} \text{tr}(Np) \quad (15)$$

$$= S_{\text{BF}+\text{particle}}^{3d}, \quad (16)$$

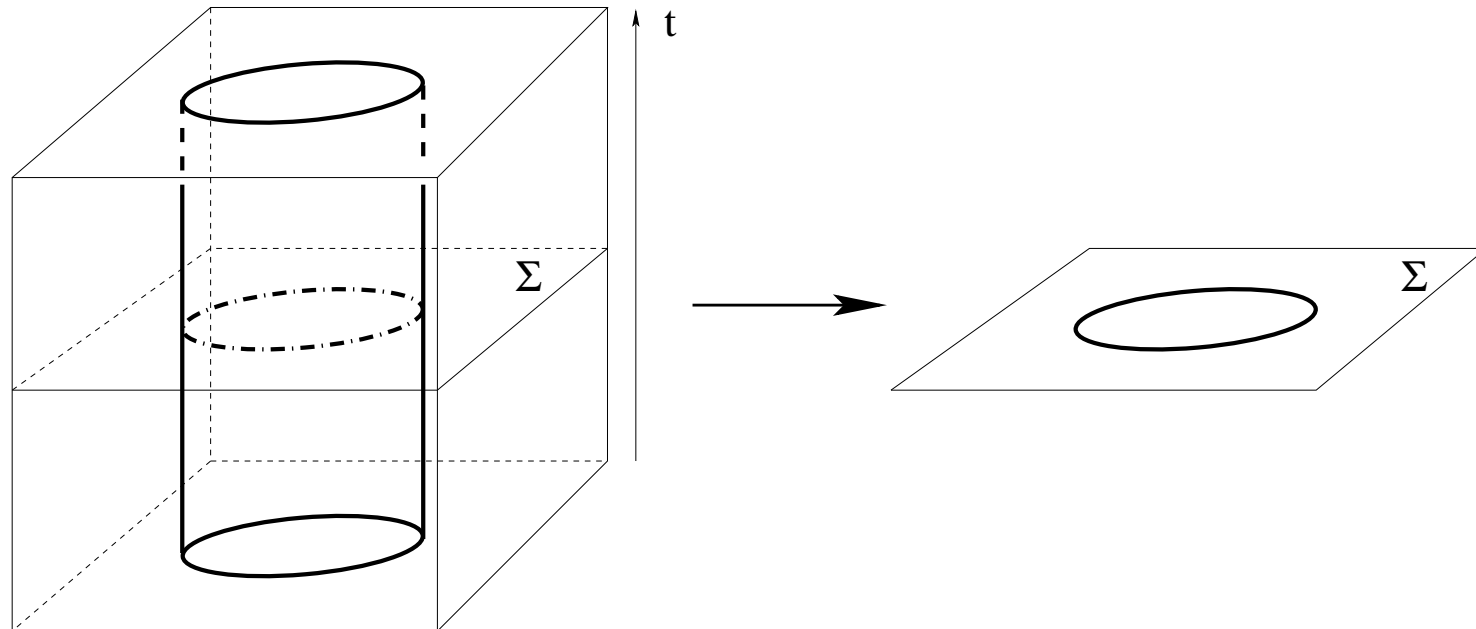
⇒ Action of 3d GR coupled to a point particle (Freidel, Louapre - 04)

Introducing the linear form  $P : \text{Cyl} \rightarrow \mathbb{C}; \Psi \mapsto \langle \Omega, \delta(\hat{C})\Psi \rangle$ , we have furthermore ( $d > 3$ )

$$P(\Omega)_{\text{BF}+(d-3)\text{-branes}}^d = \mathcal{Z}_{\text{BF}+(d-4)\text{-branes}}^{d-1} \quad (17)$$

where  $\mathcal{Z}^{d-1} \equiv$  path integral of  $(d-1)$ -dimensional BF theory coupled to branes

Quantization : physical inner product



$\Rightarrow$  Transition amplitudes of 4d BF theory coupled to strings dual to evaluations of Feynman loops coupled to 3d quantum gravity

(Barrett - 04, 05; Freidel, Livine - 05)

- **Regularization** of  $P[\Psi_\Gamma]$  - (Very!) heuristic idea
- Pick a **cellular decomposition**  $\mathcal{T}_\epsilon^* = (v^*, e^*, f^*)$  of  $\Sigma$  **adapted** to  $\Gamma$ , i.e., s.t.  $\Gamma \subset \partial f^* \subset \mathcal{T}_\epsilon^*$
- Impose  $F = 0$  ( $g_{f^*} = \mathbb{1}$ ) around **all two-cells**  $f^*$  except for the **two-cells circling the string** where  $F = p$  ( $g_{f^*} = \exp p$ )
- The **physical inner product** yields  $P[\Psi] = \lim_{\epsilon \rightarrow 0} P[\mathcal{T}_\epsilon^*; \Psi]$ , with

$$P[\mathcal{T}_\epsilon^*; \Psi] = \langle \Omega, \left[ \prod_{f^* \notin \mathcal{S}} \hat{\delta}(g_{f^*}) \prod_{f^* \in \mathcal{S}} \hat{\delta}(g_{f^*} \exp p) \right] \Psi \rangle \quad (18)$$

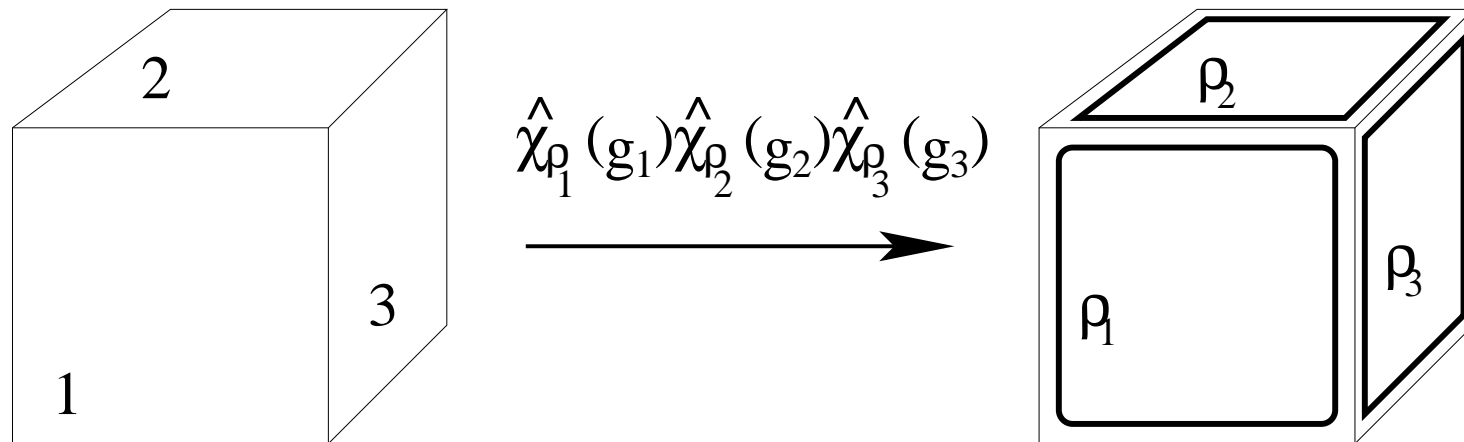
- Remark : There are in fact **many subtleties** (**reducibility** of the constraints ( $D_a C_i^a = 0$ ), presence of **open SN edges**...)
- The **full regulator**  $R_\epsilon$  is not simply the cellular complex  $\mathcal{T}_\epsilon^*$ ; it contains **other structures** (a maximal tree etc...)

Quantization : physical inner product

-Each delta function can then be given an operational meaning

→ Peter-Weyl decomposition  $\delta(g) = \sum_{\rho} \dim(\rho) \chi_{\rho}(g)$

→  $\hat{\chi}_{\rho}(g)$  self-adjoint Wilson loop operator on  $\mathcal{H}_{kin}$





### Topological invariance

- How to remove the regulator  $R_\epsilon$  ?

→ We show that

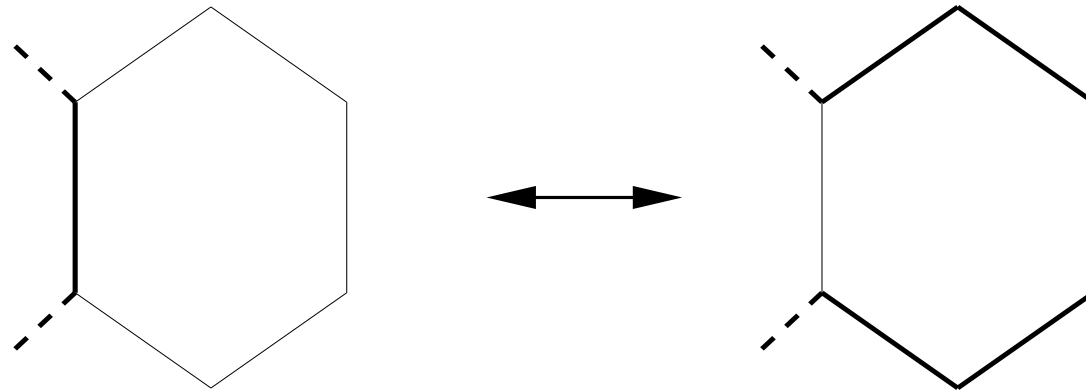
$$\forall \Psi \in \mathcal{H}_{kin}, \quad P[R_\epsilon; \Psi] = P[R'_\epsilon; \Psi] \quad (19)$$

where,

$R_\epsilon \mapsto R'_\epsilon$ , finite combination of discrete moves acting on each component of the full regulator  $R_\epsilon$  (3d Pachner moves, elementary tree moves ...)

- Furthermore, we can show the invariance of the physical inner product under elementary SN edge moves (map PL-paths into ambient isotopic PL-paths)

Topological invariance



$\Rightarrow$  Putting all this together, we can show that

$$\forall \Psi_{\Gamma} \in \mathcal{H}_{kin}, \quad P[R_{\epsilon}; \Psi_{\Gamma}] = P[[\Sigma]; \Psi_{[\Gamma]}] \quad (20)$$

where  $[\Sigma]$  and  $[\Gamma]$  denote the **equivalence classes** of topological **manifolds** and **one-complexes** up to homeomorphisms and ambient isotopy respectively

## Conclusion and Outlook

- The **road** towards a **clear physical picture** of the topological models discussed here **is open**
- The **physical inner product** between any two states can be **explicitly computed**
- The **transition amplitudes** are **topological invariants** of the canonical manifold together with the embedded string spin networks
- To do :
  - Write the **covariant model**
  - Re-express the **amplitudes** as **Feynman diagrams of a QFT** (String field theory ?)
  - Add the Plebanski term** (contact with conventional ST ?)
  - ...