

## Group Field Theories: spacetime from quantum discreteness to an amergent continuum

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- 6 conclusions



6 best formalism we have for microscopic physics and many-particle physics (particle physics, atomic physics, condensed matter,...)



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- $\begin{tabular}{ll} \hline & \phi(x) \end{tabular} \rightarrow S_{\lambda}(\psi) \end{tabular} \rightarrow Z = \int D\psi \, e^{iS_{\lambda}(\psi)} \end{tabular}$



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6 discrete regime: few quanta 
$$\rightarrow$$
  
 $\rightarrow Z = \sum_{\Gamma} \frac{\lambda^{N_{\Gamma}}}{sym\Gamma} A(\Gamma)$   
 $A(\Gamma) = \prod_{v} \int dx_{v} A(\Gamma, x_{v}) = \prod_{e} \int dp_{e} A(\Gamma, p_{e})$   
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6 continuum regime: many quanta  $\rightarrow$   $S(\psi); Z = \int D\psi e^{iS_{\lambda}(\psi)}$  or  $H(\psi); Z = \int D\psi e^{-\beta H} \Rightarrow$   $\Rightarrow S_{eff}(\phi); Z_{eff} = \int D\phi e^{iS_{eff}(\phi)}$  or  $H_{eff}(\phi); Z_{eff} = \int D\phi_{eff} e^{-\beta H_{eff}(\phi)}$ 





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  - we are not starting from scratch! ideas and results from LQG, matrix models, simplicial QG,...
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#### 6 GFTs are

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- 6 arguments of field have interpretation of geometric data (lengths, areas, gravity connection,...)



- 6 GFT action  $S(\phi)$  with  $V(\phi)$  of order D+1, and non-local pairing of field variables (combinatorics of D-simplex)
- **6** FD are fat graphs/2-complexes  $\Gamma$  topologically dual to D-dimensional simplicial complexes  $\Delta$  (discrete spacetime emerges as virtual construct)

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- Generalization of matrix models for 2d QG: more complicated combinatorics + group-theoretic data
- both geometry and topology are dynamical


# The discrete regime

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#### The Group Field Theory formalism

#### L. Freidel, hep-th/0505016; D. Oriti, gr-qc/0512103; D. Oriti, gr-qc/0607032



D

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$$\phi (g_1 g_2 g_3) \longleftrightarrow \phi(j_1 j_2 j_3) \longleftrightarrow \underbrace{\varphi_1(j_1)}_{g_3(j_3)} \underbrace{g_2(j_2)}_{g_3(j_3)} \underbrace{\varphi_1(j_1)}_{g_3(j_3)} \underbrace{g_2(j_2)}_{g_3(j_3)} \underbrace{\varphi_2(j_2)}_{g_3(j_3)} \underbrace{\varphi_2(j_$$



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<sup>6</sup> generic state  $\rightarrow$  spin network  $\simeq$  simplicial (D-1)-complex, with geometric data attached



6 field action:  $S_D(\phi, \lambda) = \frac{1}{2} \prod_{i=1,..,D} \int dg_i d\tilde{g}_i \phi(g_i) \mathcal{K}(g_i \tilde{g}_i^{-1}) \phi(\tilde{g}_i) + \frac{\lambda}{(D+1)} \prod_{i\neq j=1}^{D+1} \int dg_{ij} \phi(g_{1j}) \dots \phi(g_{D+1j}) \mathcal{V}(g_{ij} g_{ji}^{-1})$ exact choice of the  $\mathcal{K}$  and  $\mathcal{V}$  defines the model



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$$D = 3$$
):  $g \in SU(2)$ 

$$\mathcal{K}(g_i, \tilde{g}_i) = \prod_i \, \delta(g_i \tilde{g}^{-1}) \quad \mathcal{V}(g_{ij}) = \prod_{i \neq j} \delta(g_{ij} g_{ji}^{-1})$$

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o extensive study of classical dynamics carried out yet



$$Z = \int \mathcal{D}\phi \, e^{iS[\phi]} = \sum_{\Gamma} \, \frac{\lambda^{N_{\Gamma}}}{sym[\Gamma]} \, Z(\Gamma)$$



6 the quantum theory is defined by the partition function, in terms of its Feynman expansion:

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- 6 this produces: c) 2-cells, identified by strands of propagation passing through several vertices, and then closing (for closed FD), dual to (D-2)-simplices; d) 'bubbles': 3-cells bounded by the above 2-cells

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- Quantum Gravity formulated as a sum over simplicial complexes of all topologies, as interaction processes



**6** Feynman amplitudes can be written in both configuration  $(g_i)$  and momentum  $(J_i)$  space

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$$Z(\Gamma) \simeq$$
 discrete QG path integral  $ightarrow {''} \int {\cal D} g_\Delta \; e^{i \, S_\Delta(g) \; ''}$ 



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$$\langle \Psi_1 \mid \Psi_2 \rangle = \int \mathcal{D}\phi O_1 O_2 e^{iS(\phi)} = \sum_{\Gamma/\partial\Gamma = \gamma_1 \cup \gamma_2} \frac{\lambda^{N_{\Gamma}}}{sym[\Gamma]} Z(\Gamma)$$

- 6  $Z(\Gamma) \simeq \text{discrete QG path integral} \rightarrow \ '' \int \mathcal{D}g_{\Delta} \ e^{i S_{\Delta}(g)} \ ''$
- in most models,  $Z(\Gamma)$  directly related to/derived from discretization of continuum gravity action





**6** 3d Riemannian QG - D = 3, G = SU(2)

## Example



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6 same result from path integral quantization of 3d Riemannian gravity on triangulation  $\Delta$  dual to  $\Gamma$ 

$$S_{\mathcal{M}}(e,\omega) = \int_{\mathcal{M}} tr(e \wedge F(\omega)) \to S_{\Delta}(X_e, g_{e*}) = \sum_e tr(X_e G_e)$$



so that the  $g_{e*}$  play the role of discretized connection

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$$\begin{cases} j_1 & j_2 & j_3 \\ j_4 & j_5 & j_6 \end{cases} \propto_{j \to \infty} e^{i S_R(j_e)} + e^{-i S_R(j_e)}, \\ S_R = \text{Regge action, } j_e = \text{edge lengths} \end{cases}$$



#### GFTs as a general framework for (discrete) Quantum Gravity approaches

Group Field Theories: spacetime from quantum discreteness to an amergent continuum - p. 18/3



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(S. Drappeau, E. Livine, D.O., in progress)



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- $\begin{tabular}{ll} \hline & \Psi(x_1,x_2,\ldots) \rightarrow \psi(x) \rightarrow S(\phi),\ldots \end{tabular}$

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- GFT perturbative expansion at tree level (trivial topology) defines physical scalar product of LQG !!! (L. Freidel)

$$\langle \Psi_1 \mid \Psi_2 \rangle_{phys} = \sum_{\Gamma \mid tree / \partial \Gamma = \gamma_1 \cup \gamma_2} \frac{\lambda^{N_{\Gamma}}}{sym[\Gamma]} Z(\Gamma)$$

Group Field Theories: spacetime from quantum discreteness to an amergent continuum - p. 22/3





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- ▲ beyond tree level: there is much more than that in a QFT, → there is much more in GFT than in LQG....

6 GFT seems to incorporate formalism, insights and ideas from both main approaches to simplicial quantum gravity:

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Dynamical Triangulations: freeze the geometric data, then GFT gives QG as sum over triangulations

$$Z = \int \mathcal{D}\phi \, e^{iS[\phi]} \curvearrowright Z_{DT} = \sum_{\Gamma} \frac{1}{sym[\Gamma]} \, A_{\Gamma}(\lambda) \approx \int Dg \, e^{iS_{GR}(g)},$$

(too much included? need for further restriction to 1) trivial topology, 2) Causal Dynamical Triangulations? QFT meaning of CDT restrictions?)



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  - ▲ may be useful to introduce extra variables (not g-only or j-only) to reproduce  $(e, \omega)$  or  $(L_e, g_{e*})$  Group Field Theories: spacetime from quantum discreteness to an amergent continuum - p. 25/3

(D.O., T. Tlas, in preparation)



6 consider the GFT, for D = 3, both Riemannian and Lorentzian, based on the complex field  $\phi^{\alpha}(g_1, s_1; g_2, s_2; g_3, s_3) : [(G \times \mathbb{R})^{\times 3}/G] \to \mathbb{C}$ , with  $\phi^{\alpha=+1} = \phi$ ,  $\phi^{\alpha=-1} = \phi^*$  (orientation dependence)

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6 and the action  

$$S = \int_{G} dg_{i} \int_{\mathbb{R}} ds_{i} \phi^{*}(g_{i}, s_{i}) \left(\prod_{i} (i\partial_{i} + \Box_{i})\right) \phi(g_{i}, s_{i}) + \sum_{\alpha_{i}} \frac{\lambda}{4} \int dg_{ij} ds_{ij} \phi^{\alpha_{1}}(g_{1j}, s_{1j}) \dots \phi^{\alpha_{4}}(g_{4j}, s_{4j}) \prod_{ij} \delta(g_{ij}, g_{ji}) \delta(\alpha_{i} s_{ij} + \alpha_{j})$$

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(D.O., T. Tlas, in preparation)



6 consider the GFT, for D = 3, both Riemannian and Lorentzian, based on the complex field φ<sup>α</sup>(g<sub>1</sub>, s<sub>1</sub>; g<sub>2</sub>, s<sub>2</sub>; g<sub>3</sub>, s<sub>3</sub>) : [(G × ℝ)<sup>×3</sup>/G] → ℂ, with φ<sup>α=+1</sup> = φ, φ<sup>α=-1</sup> = φ<sup>\*</sup> (orientation dependence)

6 and the action  

$$S = \int_{G} dg_{i} \int_{\mathbb{R}} ds_{i} \phi^{*}(g_{i}, s_{i}) \left(\prod_{i} (i\partial_{i} + \Box_{i})\right) \phi(g_{i}, s_{i}) + \sum_{\alpha_{i}} \frac{\lambda}{4} \int dg_{ij} ds_{ij} \phi^{\alpha_{1}}(g_{1j}, s_{1j}) \dots \phi^{\alpha_{4}}(g_{4j}, s_{4j}) \prod_{ij} \delta(g_{ij}, g_{ji}) \delta(\alpha_{i} s_{ij} + \alpha_{j})$$

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 $\circ$  usual model(s) = SUL restriction:  $\prod_i (i\partial_i + \Box_i) \rightarrow \prod_i (\delta(g_i, \tilde{g}_i)\delta(s_i - \tilde{s}_i))$ 

Group Field Theories: spacetime from quantum discreteness to an amergent continuum - p. 26/3

(D.O., T. Tlas, in preparation)



$$Z_{\Gamma}(g_{e*}, M_e) = \left[\prod_{e} C(N_e) \left(\frac{1}{\sqrt{1+M_e}\Theta_e(g_{e*})}\right)^{N_e-1} \times \sum_{k=0}^{N_e-1} \frac{(N_e+k-1)!}{k!(N_e-k-1)!} \frac{(-1)^k}{(2i\sqrt{1+M_e})^k}\right] e^{i\sum_{e} \sqrt{1+M_e}\Theta_e(g_{e*})} \quad M_e > -1$$

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- **5** DT restriction (fixed geometric data):  $\phi(g_i, s_i)$  eigenstate of  $i\partial_i$ , trivial gauge invariance (D.O., in progress) Group Field Theories: spacetime from quantum discreteness to an amergent continuum – p. 27/3

# The problem of the emergence of the continuum (and of GR) from GFTs

#### Some current strategies (very much simplified!) seen from GFT

$$\langle \Psi_1 \mid \Psi_2 \rangle = \int \mathcal{D}\phi O_1(\phi) O_2(\phi) e^{iS(\phi)} = \sum_{\Gamma \mid \gamma_1 \cup \gamma_2} \frac{\lambda^{N_{\Gamma}}}{s[\Gamma]} \sum_{\{J_i\}} A(J_i)$$

<sup>6</sup> LQG/SF semiclassical states: construct kinematical spin network states such that  $\gamma \simeq \Sigma$  and  $\{J_i\} \simeq h_{\Sigma}(x)$ ; then study their dynamics using an appropriate spin foam model, in which they would appear as boundary states

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- 6 GFT translation: identify appropriate hugely populated ( $10^{24}$ ?  $10^{51}$ ?) "multi-particle states"(each with its own "momentum"), that can be approximated by a continuum, characterized by a smooth "momentum field". Study their dynamics in perturbative expansion around the vacuum, analyzing hugely complicated FD amplitudes (complexity of the FD  $\simeq$ complexity of the states)



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- 6 if continuum is many-particle physics, NO!



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  - in general, (modified) GR does not coincide with hydrodynamics, but is reproduced only in some limits
- 6 questions from CM perspective: what are the atoms of space? what is the microscopic theory? which CM system reproduces full GR?



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- 6 hypothesis: continuum is coherent many-particles physics for GFT atoms of space at very low temperature (hydrodynamic approx)



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identify GFT phases for  $T \rightarrow 0$ , large number of quanta, for different GFT models

microscopic details mostly irrelevant (e.g. exact form of Feynman amplitudes/SF), what is relevant: variables, symmetries, statistics,

general form of interaction



6 continuum spacetime = fluid phase, i.e. intertwiners/(D-1)-simplices condense or reach ground state and evolve coherently, as a fluid-continuum

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- extract effective dynamics of these hydrodynamic degrees of freedom from microscopic GFTs can use usual QFT methods for condensed matter physics + insights from analog gravity
- Iook for geometric interpretation of hydrodynamic variables and dynamics (does it give (modified) GR?)

GFTs contain topology change.....is the effective continuum theory going to be GR on a fixed topology, or rather the classical

continuum theory on superspace behind 3rd quantized gravity?

(D.O., J. Ryan, A. Youssef, in preparation)



6 consider model with kinetic term:

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for generic group G (Riemannian or Lorentzian)

6 take  $s_i$  as time variables, and  $g_i$  as 'space'variables; for each argument of field there is one time variable

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- 6 can define appropriate Poisson Brackets, conserved quantities, etc

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- In upon quantization, in order to preserve positivity of the Hamiltonian, it turns out that GFTs of this type have to be quantized using Bose statistics in Riemannian case, and Fermi statistics in Lorentzian case

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- 6 only problem: there is plenty of work still to be done! :(