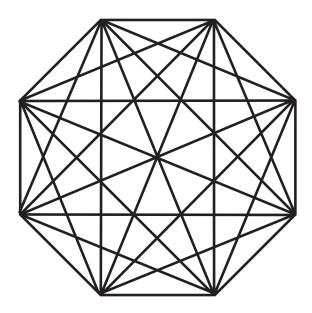
QUANTUM GRAVITY AND EMERGENT LOCALITY

Fotini Markopoulou Perimeter Institute QUANTUM GRAVITY AND EMERGENT LOCALITY QUANTUM GRAPHITY: a background independent condensed matter model of emergent space

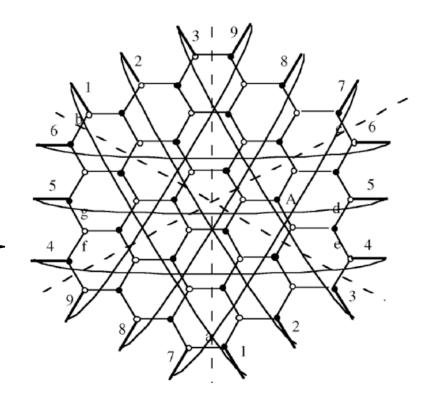
> Fotini Markopoulou Perimeter Institute

T. Konopka, FM & L.Smolin, hep-th/0611197 also with: O. Dreyer, S. Severini

Preview



geometrogenesis phase transition



 $\mathsf{High-}\,T$

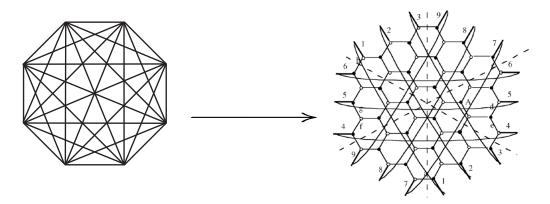
- Permutation symmetry
- No locality (no space)
- Relational

Low-T ground state

- Translations
- Local
- Relational

Outline.

- 1. Motivation
- 2. Quantum Graphity: A Model of emergent space (time)

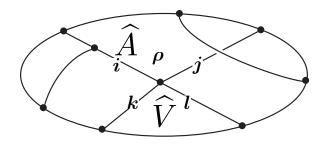


- 3. Could we observe this? Early universe application
- 4. Discussion

I. Emergence of near-flat geometry in low-energy limit of background independent quantum gravity: long-standing challenge. cf. Ambjorn talk

I. Emergence of near-flat geometry in low-energy limit of background independent quantum gravity: long-standing challenge. cf. Ambjorn talk

- 2. Locality: underutilized.
 - combinatorial quantum geometry classical geometry



 \widehat{A}, \widehat{V} : quantum geometry spin foam vertex properties $\left. \right\}$ local statements

I. Emergence of near-flat geometry in low-energy limit of background independent quantum gravity: long-standing challenge. cf. Ambjorn talk

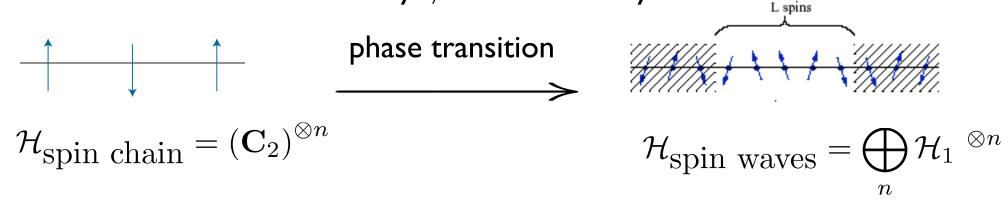
- 2. Locality: underutilized.
 - combinatorial quantum geometry classical geometry

$$\widehat{A}, \widehat{V}: \text{quantum geometry} \text{spin foam vertex properties} \right\} \text{ local statements}$$

$$N \text{ nodes} \Rightarrow N^2 \text{ one-edge perturbations} \right\} \text{ locality is unstable}$$

$$\sum_i |S_i\rangle \Rightarrow \text{ locality?} \qquad \text{FM&L.Smolin, gr-qc/0702044}$$

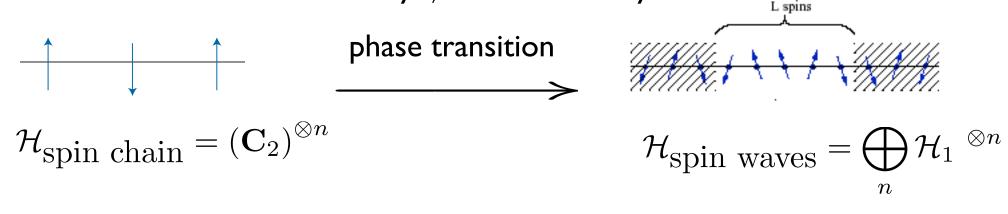
- 2. Locality, cont.
 - Phase transition + background independence \Rightarrow micro-locality \neq macro-locality



background independence means dynamics determines locality.

O.Dreyer, hep-th/0409048 FM&D.Kribs, gr-qc/0510052 FM, gr-qc/0703097

- 2. Locality, cont.
 - Phase transition + background independence \Rightarrow micro-locality \neq macro-locality



background independence means dynamics determines locality.

O.Dreyer, hep-th/0409048 FM&D.Kribs, gr-qc/0510052 FM, gr-qc/0703097

Lots of other hints: AdS/CFT, black holes, causets,...

3. Background independence vs emergent gravity

Geometry and gravity are only classical, emergent concepts.

- condensed matter approaches
 (e.g.Volovik)
- string theory

Background independence: dynamical nature of spacetime geometry

What is the role of the lattice/geometry and the symmetries of the lattice?

3. Background independence vs emergent gravity

Geometry and gravity are only classical, emergent concepts.

- condensed matter approaches (e.g.Volovik)
- string theory

Background independence: dynamical nature of spacetime geometry

What is the role of the lattice/geometry and the symmetries of the lattice?

Background independence is normally implemented by a superposition

of quantum geometries: - loop quantum gravity

- dynamical triangulations
- causal sets

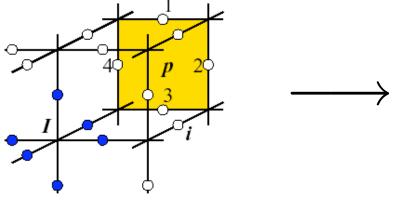
Reconcile:

Background independence: no fundamental geometric degrees of freedom

• Today: background independent condensed matter models. Dreyer 04, 06, Lloyd 05, FM&Kribs 05, FM07.

4. Matter vs Geometry.

New results on emergent matter are suitable for a background independent treatment:



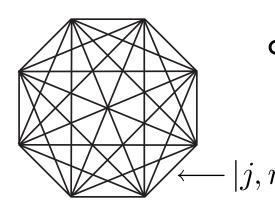
M.Levin & X-G.Wen, hep-th/0507118

gauge theory -like excitations and fermions

5. No fundamental locality \implies large-scale observational consequences of quantum gravity.

The model: kinematics.

Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges.



one-edge state space:
$$\mathcal{H}_1 = \operatorname{span}|j, m\rangle$$

 $\bullet \quad |0, 0\rangle \quad \text{edge off}$
 $\bullet \quad |1, -1\rangle \quad |1, +1\rangle \quad edge on$
 $\bullet \quad |1, 0\rangle \quad edge on$

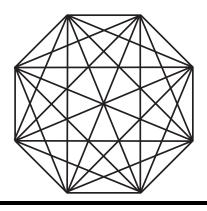
total state space:
$$\mathcal{H} = \mathcal{H}_1^{\otimes \frac{n(n-1)}{2}}$$

$$J|j,m\rangle = j|j,m\rangle \qquad M^+|1,0\rangle = |1,1\rangle$$

$$M|j,m\rangle = m|j,m\rangle \qquad M^-|1,0\rangle = |1,-1\rangle$$

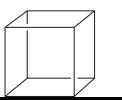
$$J|0,0\rangle = M|0,0\rangle = M^{\pm}|0,0\rangle = 0$$

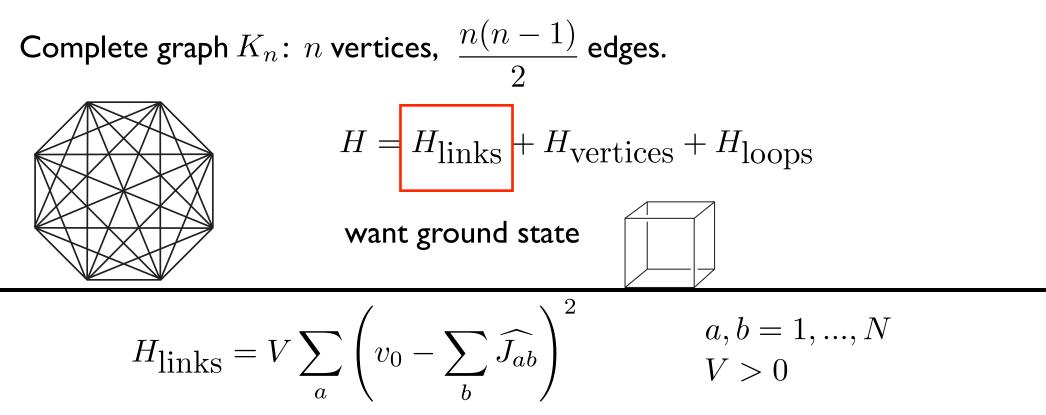
Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges.



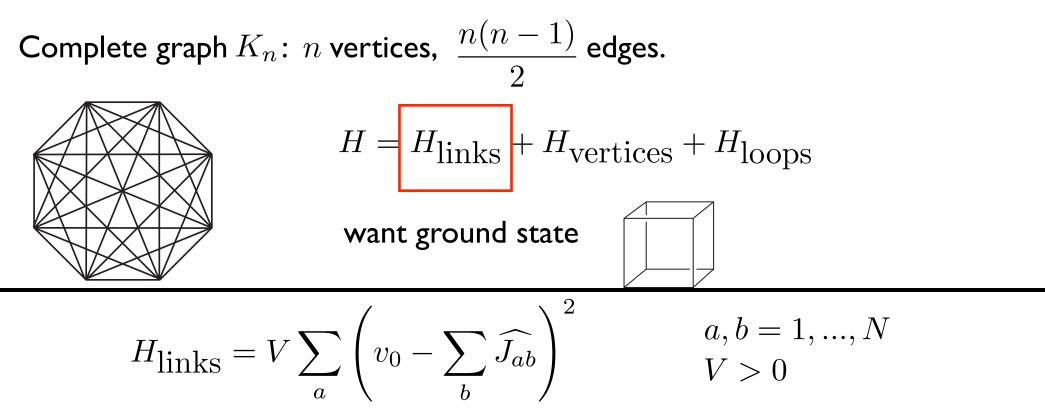
$$H = H_{\text{links}} + H_{\text{vertices}} + H_{\text{loops}}$$

want ground state

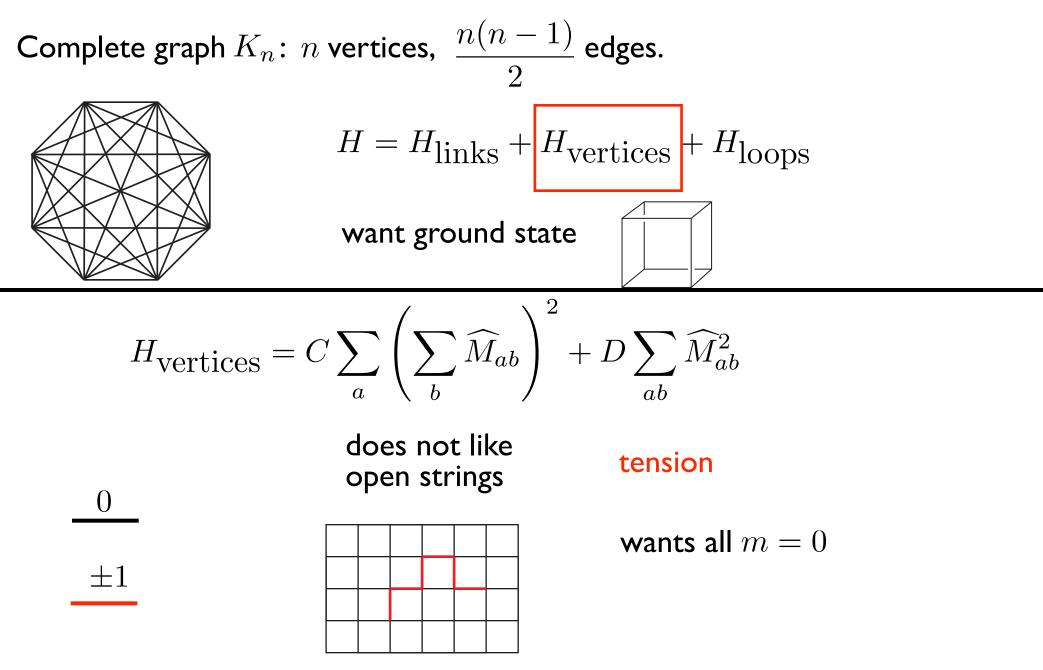




 H_{links} turns off edges and has minimum when all edges have degree v_0 . E.g., for $v_0 = 3$,



 H_{links} turns off edges and has minimum when all edges have degree v_0 . E.g., for $v_0 = 3$, m values: strings



i=1

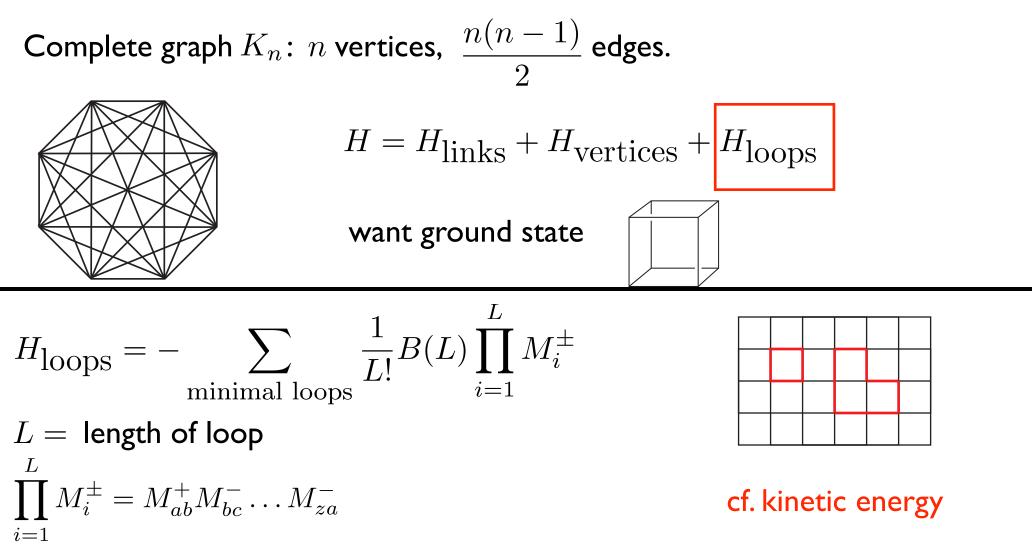
Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges. $H = H_{\text{links}} + H_{\text{vertices}} + H_{\text{loops}}$ want ground state $H_{\text{loops}} = - \sum_{i \text{ simultaneous}} \frac{1}{L!} B(L) \prod_{i=1}^{L} M_i^{\pm}$ L = length of loop $\prod M_i^{\pm} = M_{ab}^+ M_{bc}^- \dots M_{za}^$ cf. kinetic energy

i=1

Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges. $H = H_{\text{links}} + H_{\text{vertices}} + H_{\text{loops}}$ want ground state $H_{\text{loops}} = - \sum_{i \text{ simultaneous}} \frac{1}{L!} B(L) \prod_{i=1}^{L} M_i^{\pm}$ L = length of loop $\prod M_i^{\pm} = M_{ab}^+ M_{bc}^- \dots M_{za}^$ cf. kinetic energy

i=1

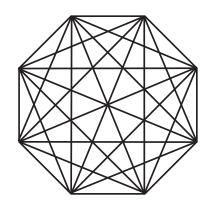
Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges. $H = H_{\text{links}} + H_{\text{vertices}} + H_{\text{loops}}$ want ground state $H_{\text{loops}} = - \sum_{i \text{ simultaneous}} \frac{1}{L!} B(L) \prod_{i=1}^{L} M_i^{\pm}$ L = length of loop $\prod M_i^{\pm} = M_{ab}^+ M_{bc}^- \dots M_{za}^$ cf. kinetic energy



Let $B(L) = B_0 B^L$. There is a preferred loop length L_* :

$$\frac{B^{L_*}}{L_*!} > \frac{B^{L'}}{L'!} \qquad \forall \ L' \neq L_*$$

Complete graph K_n : *n* vertices, $\frac{n(n-1)}{2}$ edges.



$$H = H_{\text{links}} + H_{\text{vertices}} + H_{\text{loops}}$$

$$H_{\text{links}} = V \sum_{a} \left(v_0 - \sum_{b} \widehat{J_{ab}} \right)^2$$
$$H_{\text{vertices}} = C \sum_{a} \left(\sum_{b} \widehat{M}_{ab} \right)^2 + D \sum_{ab} \widehat{M}_{ab}^2$$
$$H_{\text{loops}} = -\sum_{\text{minimal loops}} \frac{1}{L!} B(L) \prod_{i=1}^L M_i^{\pm}$$

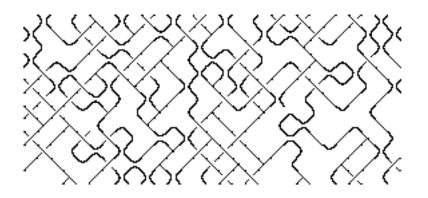
 $V \gg C, D~~$ erases edges to give lattice with extension $C \gg D~~B_0 \gg D~~$ makes lattice regular and gives macro-matter

Matter.

 $H_{\rm vertices}$ and $H_{\rm loops}$ are generalizations of the rotor model of string-net condensation of Levin and Wen to a dynamical lattice. N.Levin & X-G.Wen, hep-th/0507118

- tension \gg K.E.: ground state has almost no strings
- tension \ll K.E.: ground state is a superposition of many closed strings,

i.e., string-net condensed:



high-energy charged excitations: open strings

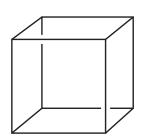


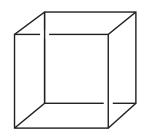
emergent U(1) gauge theory.

Ground state geometry plus matter. $N \gg 1, V \gg 1, C \gg D, B_0 \gg D$

In our model: the balance between tension and K.E. controls both the loop size when the lattice is dynamical and gives the matter.

Local lattice: large average distance $\langle d_{ij} \rangle$. Small vertex degree $\sim v_0$ Loops of average size L_* . Possibly also regular?

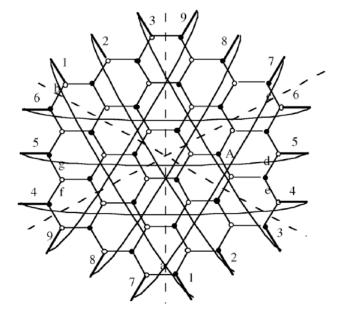




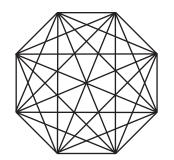
For $v_0 = 3, L_* = 6$ number of minimal loops is maximized by the honeycomb lattice:

T.Konopka, FM & S.Severini

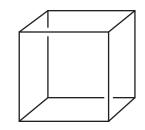
Plus string condensation on the ground state lattice, as in string nets.



Summary of quantum graphity.



geometrogenesis phase transition



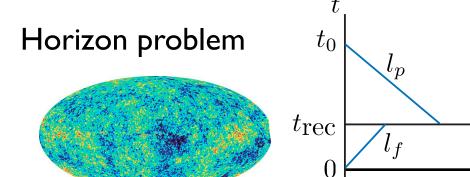
High-T

- Permutation symmetry
- No locality
- Relational
- $\langle d_{ij} \rangle = 1$
- $\sim \infty$ -dimensional
- no subsystems
- external time
- micro-matter

Low-T

- Translations
- Local
- Relational
- $\langle d_{ij} \rangle$ large
- low-dimensional
- subsystems
- external and internal time
- macro-matter and dynamical geometry

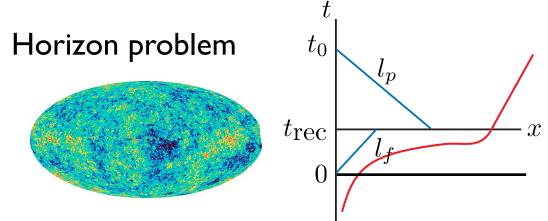
Model has 4 couplings (V, C, D, B_0) and 3 parameters (N, L_*, v_0) . Some obey generic conditions: $N \gg 1$, $V \gg 1$, $C \gg D$, $B_0 \gg D$ Some have to be adjusted: $v_0 = 3, 4$, $L_* = 4, 6$



 $l_p \gg l_f$

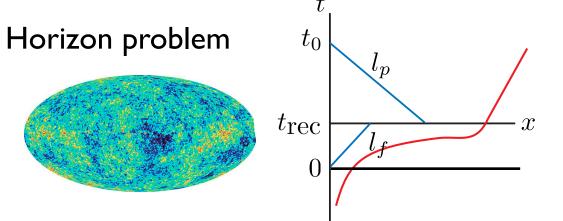
x

This is an extrapolation of our causal and local structure to the initial time



 $l_p \gg l_f$

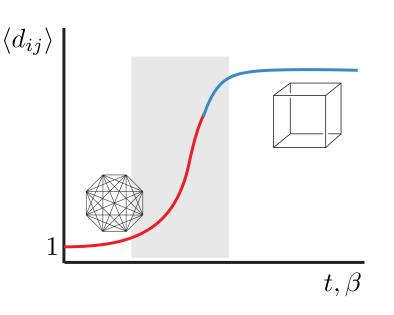
This is an extrapolation of our causal and local structure to the initial time inflation

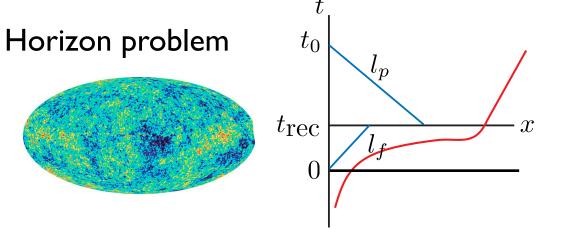


 $l_p \gg l_f$

This is an extrapolation of our causal and local structure to the initial time inflation

In our model:

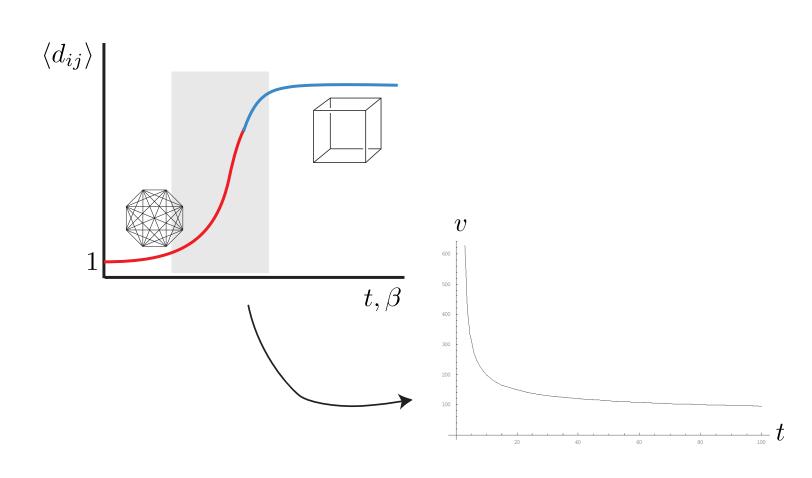


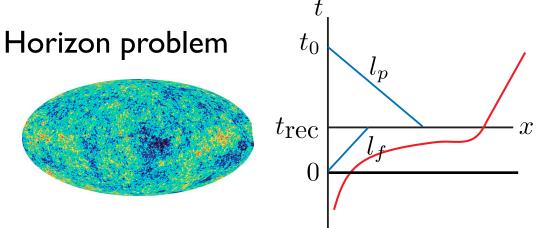


 $l_p \gg l_f$

This is an extrapolation of our causal and local structure to the initial time inflation

In our model:

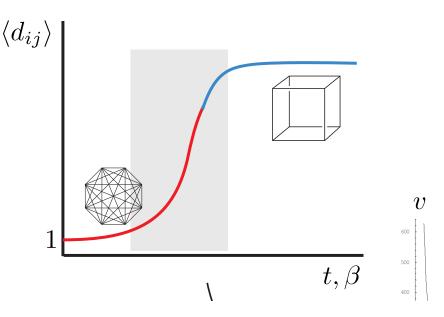




 $l_p \gg l_f$

This is an extrapolation of our causal and local structure to the initial time inflation

In our model:



Are the CMB correlations a consequence of pre-geometry? Compare, e.g., to the phenomenology of $\mathcal{O}(l_{\text{Pl}})$ discretization. _____ A geometrogenesis scenario has the potential of large-scale effects. _____ t

Comments.

I. Time, temperature? Emergent space or emergent spacetime?

• Time-reparametrization invariance as the noiseless sector of a theory with time: T. Konopka and FM, gr-qc/0601028.

• Does the temperature imply a bath?

Internal vs external time and geometry.
 Internal: subsystems observing subsystems. Our emergent space gives emergent spacectime internally.

2. The terms that give rise to matter are also responsible for the geometry of the lattice: we don't just erase edges, we also order them. This is a new use for unification.

To do.

- Understand the ground state
- Graph theory analysis of the terms in the Hamiltonian
- Cosmology: perturbations
- Effective disordered locality description
- Understand the role of the temperature

