

## Effective Constraints of Loop Quantum Gravity Mikhail Kagan

Institute for Gravitational Physics and Geometry, Pennsylvania State University

in collaboration with

M. Bojowald, G. Hossain, (IGPG, Penn State) H.H.Hernandez, A. Skirzewski (Max-Planck-Institut für Gravitationsphysik, Albert-Einstein-Institut, Potsdam, Germany



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- 2. Effective approximation. Overview.
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- 5. Summary

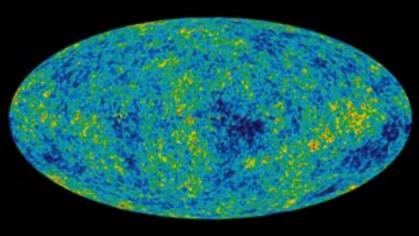


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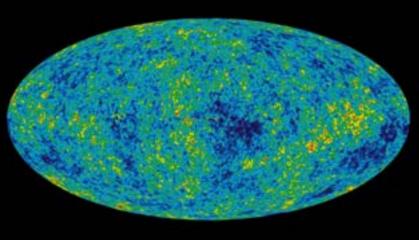
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Evolution of inhomogeneities is expected to explain cosmological structure formation and lead to observable results.



Effective approximation allows to extract predictions of the underlying quantum theory without going into consideration of quantum states.

# **Effective Approximation.** Strategy.

**Classical Theory** 

Classical Constraints & { , }<sub>PB</sub>

Quantization

Quantum Operators & [,]

**Effective Theory** 

Quantum variables:  $(q,p) \rightarrow G_q^{a,n} = \langle (\hat{q} - \langle q \rangle)^{n-a} (\hat{p} - \langle p \rangle)^a \rangle$ expectation values spreads, deformations, etc.

Truncation

classically well behaved expressions

> Classical Expressions

classically diverging expressions

**Expectation Values** 

**Classical Expressions** × **Correction Functions** 

# **Effective Approximation.** Summary. **Classical Constraints & Poisson Algebra** quantization **Constraint Operators & Commutation Relations** effective ] approximation **Effective Constraints & Effective Poisson Algebra** (differs from classical Poisson Algebra) **Anomaly Issue Effective Equations of Motion** (Bojowald, Hernandez, MK, Singh, Skirzewski Phys. Rev. D, 74, 123512, 2006; Phys. Rev. Lett. 98, 031301, 2007 for scalar mode in longitudinal gauge)

# **Effective Approximation.** Summary. **Classical Constraints & Poisson Algebra** quantization **Constraint Operators & Commutation Relations** effective ] approximation **Effective Constraints & Effective Poisson Algebra** (differs from classical Poisson Algebra) **Anomaly Issue Effective Equations of Motion** (Bojowald, Hernandez, MK, Singh, Skirzewski Phys. Rev. D, 74, 123512, 2006; Phys. Rev. Lett. 98, 031301, 2007 for scalar mode in longitudinal gauge)

# **Anomalies.** Source of Corrections.

**Basic Variables** 

Densitized triad $E_i^a$ Ashtekar connection $A_a^i = \Gamma_a^i + \gamma K_a^i$ Scalar field $\varphi$ Field momentum $\pi$ 

### **Diffeomorphism Constraint**

intact

Hamiltonian constraint

$$H_g[N] = \frac{1}{16\pi G} \int_{\Sigma} d^3x \frac{N}{\sqrt{|\det E|}} \left( \epsilon_{ijk} F^i_{ab} E^a_j E^b_k - 2(1+\gamma^2) K^i_a K^j_b E^{[a}_i E^{b]}_j \right)$$
$$H_m[N] = \int d^3x N \left[ \frac{\pi^2}{2\sqrt{|\det E|}} + \frac{E^a_i E^b_i}{2\sqrt{|\det E|}} \partial_a \phi \partial_b \phi + \sqrt{|\det E|} U(\phi) \right]$$

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$$\alpha(E) \qquad D(E) \qquad \sigma(E)_{0}$$

## **Anomalies.** Restrictions on Correction Functions.

**Corrected constraint algebra** 

 $H^Q = H_g^Q + H_m^Q, \quad D^Q \equiv D = D_g + D_m$ 

$$\{H^{Q}[N], D[N^{a}]\} = -H^{Q}[\tilde{N}] + \int d^{3}x N_{c} \partial N^{j} \left(E_{i}^{c} \delta_{j}^{a} - E_{j}^{c} \delta_{i}^{a}\right) \\ \times \left[\frac{\partial \alpha}{\partial E_{i}^{a}}(...) + \frac{\partial D}{\partial E_{i}^{a}}(...) + \frac{\partial \sigma}{\partial E_{i}^{a}}(...)\right]$$

 $\{H^{Q}[N_{1}], H^{Q}[N_{2}]\} = D_{g}[\alpha^{2}(N_{1}\partial^{a}N_{2} - N_{2}\partial^{a}N_{1})] + D_{m}[D\sigma(N_{1}\partial^{a}N_{2} - N_{2}\partial^{a}N_{1})]$ 

## **Anomalies.** Restrictions on Correction Functions.

**Anomaly free conditions:** 

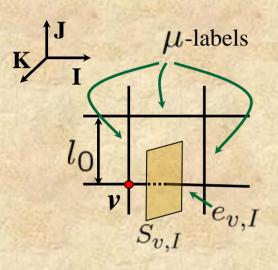
 $\alpha, D, \sigma = f(\mathcal{E}(E_i^a))$ 

 $\left(E_i^c \delta_j^a - E_j^c \delta_i^a\right) \frac{\partial \mathcal{E}}{\partial E_i^a} = 0$ 

 $\alpha^2 = D\sigma$ 

*f*=arbitrary function

### Effective Constraints. Lattice formulation. (scalar mode/longitudinal gauge, Bojowald, Hernandez, MK, Skirzewski, 2007)



Fluxes  $E_i^I = \tilde{p}^I(x)\delta_i^I$   $\downarrow$   $p_{v,I} \approx l_0^2 \tilde{p}^I(v)$ 

(integrated over  $S_{v,I}$ )

Holonomies  $K_I^i = \tilde{k}_I(x)\delta_I^i$   $\eta_{v,I} \approx \exp\left(il_0\tilde{k}_I(v)\right)$ (integrated over  $e_{vI}$ )

#### **Basic operators:**

$$\hat{p}_{v,I}|\ldots,\mu_{v',J},\ldots\rangle = 4\pi\gamma\ell_{\mathsf{P}}^{2}\mu_{v,I}|\ldots,\mu_{v',J}+1,\ldots\rangle.$$
$$\hat{\eta}_{v,I}|\ldots,\mu_{v',J},\ldots\rangle = |\ldots,\mu_{v,I}+1,\ldots\rangle.$$

# **Effective Constraints.** Hamiltonian.

#### Curvature

 $F(A) = dA + AA = (\gamma dK + \gamma^2 KK) + (d\Gamma + \Gamma\Gamma) + \gamma(\Gamma K + K\Gamma)$ 

$$\frac{FEE - KKEE}{\sqrt{|\det E|}} = \begin{bmatrix} (\gamma dK - KK) + (d\Gamma + \Gamma\Gamma) \end{bmatrix} \frac{EE}{\sqrt{|\det E|}} \\ H_{\mathcal{K}} = \sum_{v} H_{\mathcal{K},v} \qquad H_{\Gamma} \end{bmatrix}$$

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Hamiltonian

$$\hat{H}_{\mathcal{K},\mathbf{v}} = \frac{-N(\mathbf{v})}{64\pi\gamma^2 G} \sum_{IJK} \sum_{\sigma_I \in \{\pm 1\}} \{ [s_{\mathbf{v},\sigma_I I,\sigma_J J}^- s_{\mathbf{v},\sigma_J J} c_{\mathbf{v}+\sigma_I I,\sigma_J J} \}$$

 $+s_{\mathbf{v},\sigma_{I}I,\sigma_{J}J}^{+}c_{\mathbf{v},\sigma_{J}J}s_{\mathbf{v}+\sigma_{I}\mathbf{I},\sigma_{J}J}]\hat{B}_{\mathbf{v},\sigma_{K}K}\}$ 

$$c_{\mathbf{v},I} = \cos(\frac{\gamma}{2}k_{I}(v)), \quad s_{\mathbf{v},I} = \sin(\frac{\gamma}{2}k_{I}(v)),$$
$$s_{\mathbf{v},\sigma_{I}I,\sigma_{J}J}^{\pm} := \sin\left(\frac{\gamma}{2}(k_{\sigma_{I}I}(v) \pm k_{\sigma_{I}I}(v + \sigma_{J}J))\right)$$

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$$4\sin(\gamma k_I/2)\cos(\gamma k_I/2)\sin(\gamma k_J/2)\cos(\gamma k_J/2)$$

+higher curvature corrections

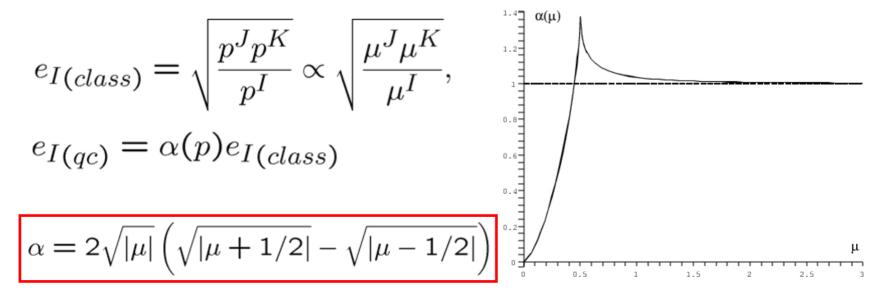
$$c_{\mathbf{v},I} = \cos(\frac{\gamma}{2}k_{I}(v)), \quad s_{\mathbf{v},I} = \sin(\frac{\gamma}{2}k_{I}(v))$$
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$$\begin{split} e_{I}^{i} &= \frac{1}{2} \epsilon^{ijk} \epsilon_{IJK} \frac{E_{j}^{J} E_{k}^{K}}{\sqrt{|\det E|}} \propto \{A_{I}^{i}, \int \sqrt{|\det E|} d^{3}x\} \longrightarrow \operatorname{tr}\left(\tau^{i}h_{v,I}[h_{v,I}^{-1}, \hat{V}]\right) \\ e_{I}^{i} &\equiv e_{I} \delta_{I}^{i} \quad \hat{e_{I}} := \hat{B}_{v,I} = \frac{1}{2\pi i \gamma \ell_{\mathsf{P}}^{2}} \operatorname{tr}\left(\tau^{i}h_{v,I}[h_{v,I}^{-1}, \hat{V}]\right) \end{split}$$

$$\hat{B}_{\mathbf{v},I}|\dots,\mu_{\mathbf{v},I},\dots\rangle := \left(4\pi\gamma\ell_{\mathsf{P}}^{2}\right)^{1/2} 2\sqrt{|\mu_{\mathbf{v},J}||\mu_{\mathbf{v},K}|} \left(\sqrt{|\mu_{\mathbf{v},I}+1/2|} - \sqrt{|\mu_{\mathbf{v},I}-1/2|}\right)|\dots,\mu_{\mathbf{v},I},\dots\rangle$$

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### Generalization

 $\frac{1}{\sqrt{|\det E|}} = \frac{(\det e)^k}{\sqrt{|\det E|^{(k+1)/2}}} \propto \left(\epsilon^{IJK}\epsilon_{ijk}\{A_I^i, V^r\}\{A_J^j, V^r\}\{A_K^k, V^r\}\right)$ <br/>for  $k \ge 1$  and  $r_k = \frac{2k-1}{3k} \ge \frac{1}{3}$ 

$$\{A_a^i, V^r\} = 4\pi\gamma G \, r V^{r-1} e_a^i$$

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#### Higher *j*-representations

$$\alpha^{(r,j)} = \frac{6}{rj(j+1)(2j+1)} |\mu|^{1-\frac{r}{2}} \sum_{m=-j}^{j} m |\mu+m|^{r/2} \text{ for large } j$$
$$\alpha^{(r,j)} = \frac{6\tilde{\mu}^{1-\frac{r}{2}}}{r(r+2)(r+4)} [(\tilde{\mu} + 1)^{\frac{r}{2}+1}(r+2-2\tilde{\mu}) + \operatorname{sgn}(\tilde{\mu}-1)|\tilde{\mu}-1|^{\frac{r}{2}+1}(r+2+2\tilde{\mu})]$$

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$$+ \text{sgn}(\tilde{\mu}-1)|\tilde{\mu}-1|^{\frac{r}{2}+1}(r+2+2\tilde{\mu})]$$

$$a^{(r,j)}_{0} a^{(r,j)}_{0} a^{(r,j)}_{0} = \frac{1+\frac{1}{32\tilde{\mu}^{2}}}{\frac{r-1}{3}} a^{(r,j)}_{0} a^{(r,j)}_{0} = \frac{1+\frac{1}{32\tilde{\mu}^{2}}}{\frac{r-2}{3}} a^{(r-2)}(r-4)\frac{4(3j^{2}+3j-1)}{5}, \quad \tilde{\mu} \to \infty$$

$$\alpha^{(r,j)} \approx (2\tilde{\mu})^{2-\frac{r}{2}}, \quad \tilde{\mu} \to 0, \qquad \tilde{\mu} := \mu/j$$



1. Can check anomalies order by order. Do not have to work with full setting.

2. Indications that ambiguities are restricted.

3. There is a consistent set of corrected constraints which are first class.

4. Cosmology: can formulate equations of motion in terms of gauge invariant variables.